# **Learning from WIMPs**

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1 Introduction



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2 Learning about the early Universe



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- 3 Learning about our galaxy

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- 4 Learning about WIMPs

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- $\Omega$ : Mass density in units of critical density;  $\Omega = 1$  means flat Universe.
- *h*: Scaled Hubble constant. Observation:  $h = 0.72 \pm 0.07$  (?)

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- Models of structure formation, X ray temperature of clusters of galaxies, ...
- Cosmic Microwave Background anisotropies (WMAP) imply  $\Omega_{\rm DM} h^2 = 0.105^{+0.007}_{-0.013}$  Spergel et al., astro-ph/0603449

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- In standard cosmology, roughly weak cross section automatically gives roughly right relic density for thermal WIMPs! (On logarithmic scale)
- Roughly weak interactions may allow both *direct* and *indirect* detection of WIMPs

# **WIMP production**

Let  $\chi$  be a generic DM particle,  $n_{\chi}$  its number density (unit: GeV<sup>3</sup>). Assume  $\chi = \overline{\chi}$ , i.e.  $\chi\chi \leftrightarrow$ SM particles is possible, but single production of  $\chi$  is forbidden by some symmetry.

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Evolution of  $n_{\chi}$  determined by Boltzmann equation:

$$\frac{dn_{\chi}}{dt} + 3Hn_{\chi} = -\langle \sigma_{\rm ann} v \rangle \left( n_{\chi}^2 - n_{\chi, \rm eq}^2 \right)$$

 $H = \dot{R}/R$ : Hubble parameter  $\langle \dots \rangle$ : Thermal averaging  $\sigma_{\rm ann} = \sigma(\chi \chi \to {\rm SM \ particles})$ v: relative velocity between  $\chi$ 's in their cms  $n_{\chi,\,{\rm eq}}: \chi$  density in full equilibrium

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Gives

$$\Omega_{\chi} h^2 \propto \frac{1}{\langle v \sigma_{\rm ann} \rangle} \sim 0.1 \text{ for } \sigma_{\rm ann} \sim \mathsf{pb}$$

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Can we test these assumptions, if  $\Omega_{\chi}$  and "all" particle physics properties of  $\chi$  are known?

MD, Imminiyaz, Kakizaki, hep-ph/0603165

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 $\sigma_{\rm ann} = a + bv^2 + \mathcal{O}(v^4) \Longrightarrow \langle \sigma_{\rm ann} v \rangle = a + 6b/x$ 

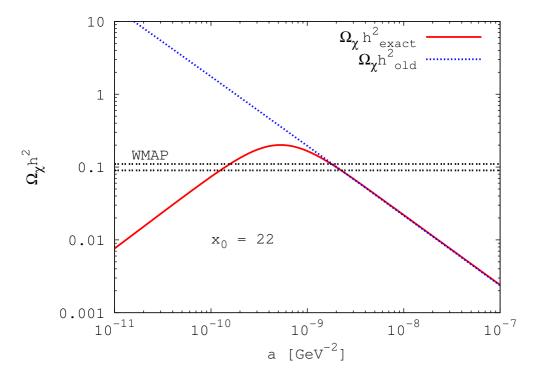
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Using explicit form of H,  $Y_{\chi,eq}$ , Boltzmann eq. becomes  $\frac{dY_{\chi}}{dx} = -f\left(a + \frac{6b}{x}\right)x^{-2}\left(Y_{\chi}^2 - cx^3e^{-2x}\right).$   $f = 1.32 \ m_{\chi}M_{\rm Pl}\sqrt{g_*}, \ c = 0.0210 \ g_{\chi}^2/g_*^2$ 

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$$Y_0(x \to \infty) = fc \left[\frac{a}{2} x_R e^{-2x_R} + \left(\frac{a}{4} + 3b\right) e^{-2x_R}\right].$$

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For intermediate temperatures,  $T_0 \lesssim T_F$ : Define 1st–order solution

$$Y_1 = Y_0 + \delta \,.$$

 $\delta < 0$  describes pure annihilation:

$$\frac{d\delta}{dx} = -f\left(a + \frac{6b}{x}\right)\frac{Y_0(x)^2}{x^2}$$

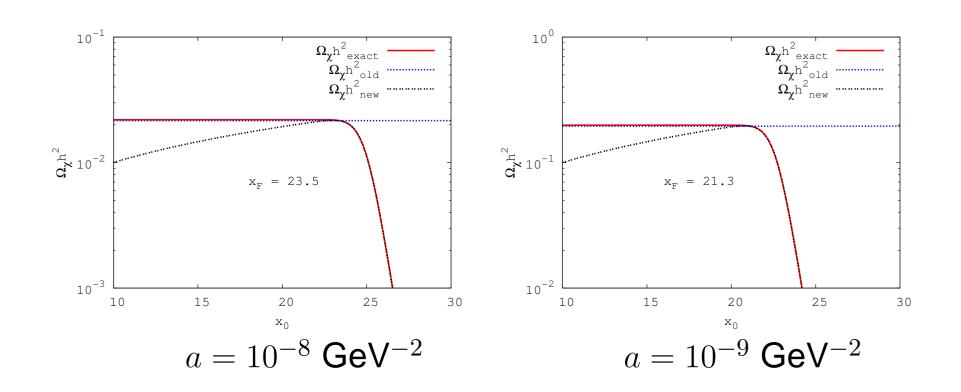
 $\delta(x)$  can be calculated analytically:  $\delta \propto \sigma_{\rm ann}^3$ 

Get good results for  $\Omega_{\chi}h^2$  for all  $T_0 \leq T_F$  through "resummation":

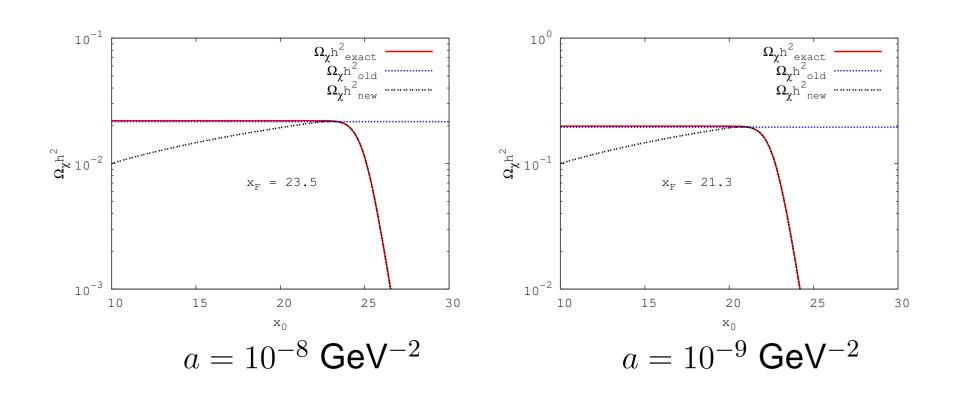
$$Y_1 = Y_0 \left( 1 + \frac{\delta}{Y_0} \right) \simeq \frac{Y_0}{1 - \delta/Y_0} \equiv Y_{1,r}$$
$$Y_{1,r} \propto 1/\sigma_{\text{ann}} \text{ for } |\delta| \gg Y_0$$

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#### **Numerical comparison:** b = 0

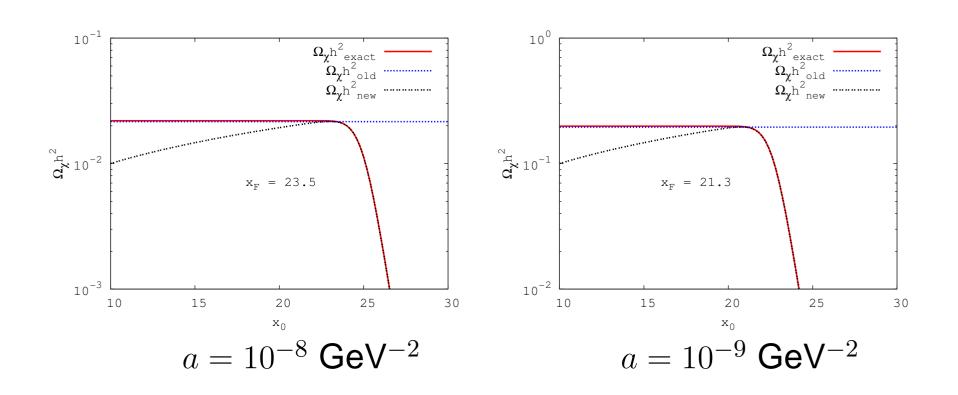


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Can extend validity of new solution to all T, including  $T \gg T_0$ , by using  $\Omega_{\chi}(T_{\text{max}})$  if  $T_0 > T_{\text{max}} \simeq T_F$ 

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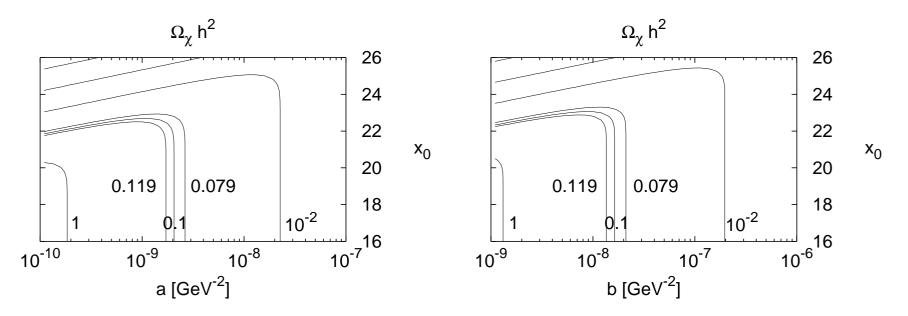


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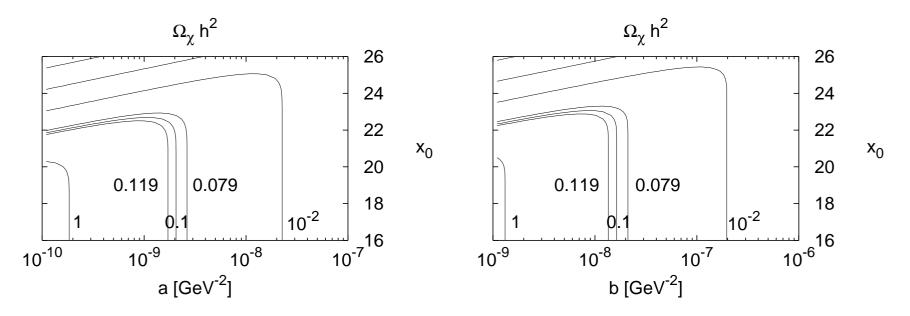
Note:  $\Omega_{\chi}(T_0) \leq \Omega_{\chi}(T_0 \gg T_F)$ 

<sup>–</sup> MD, Imminniyaz, Kakizaki, arXiv:0704.1590 [hep-ph] If  $n_{\chi}(T_0) = 0$ , demanding  $\Omega_{\chi}h^2 \simeq 0.1$  imposes lower bound on  $T_0$ :

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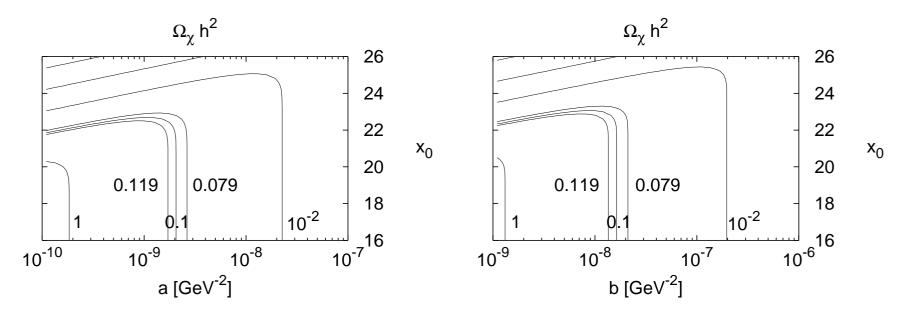
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 $\implies T_0 \ge \frac{m_{\chi}}{23}$  Holds independent of  $\sigma_{ann}!$ If  $T_0 \simeq m_{\chi}/22$ : Get right  $\Omega_{\chi}h^2$  for wide range of cross sections!

#### Assumptions

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$$A(z) = H_{\rm st}(z)/H(z)\,,\ z = T/m_{\chi}$$

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• Around decoupling:  $z \ll 1 \Longrightarrow$  use Taylor expansion

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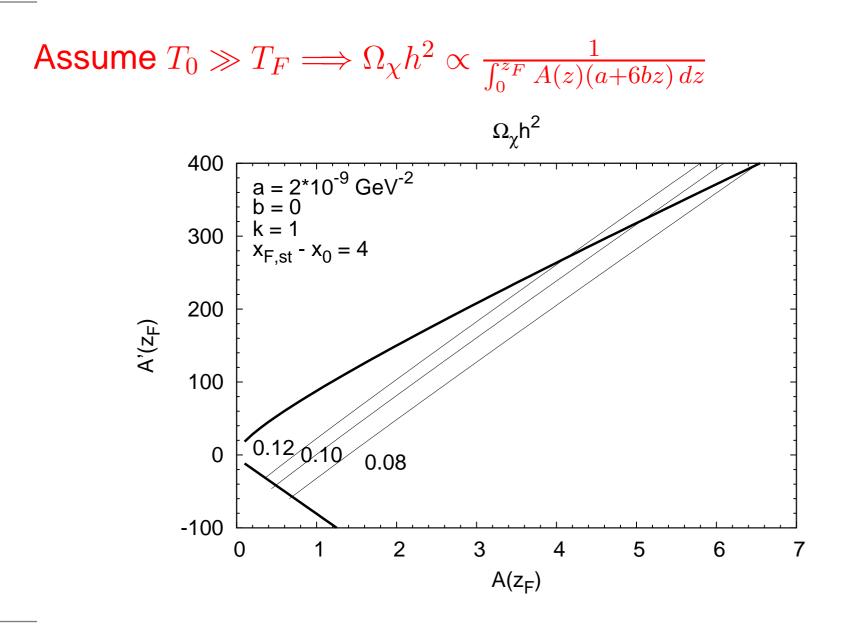
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• Successful BBN  $\implies k \equiv A(z \rightarrow 0) = 1.0 \pm 0.2$ 

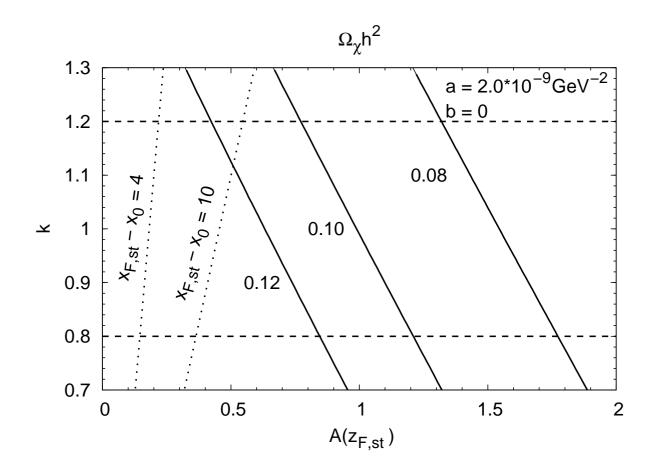
### **Constraining** H(T) (cont.d)

Assume  $T_0 \gg T_F \Longrightarrow \Omega_{\chi} h^2 \propto \frac{1}{\int_0^{z_F} A(z)(a+6bz) dz}$ 

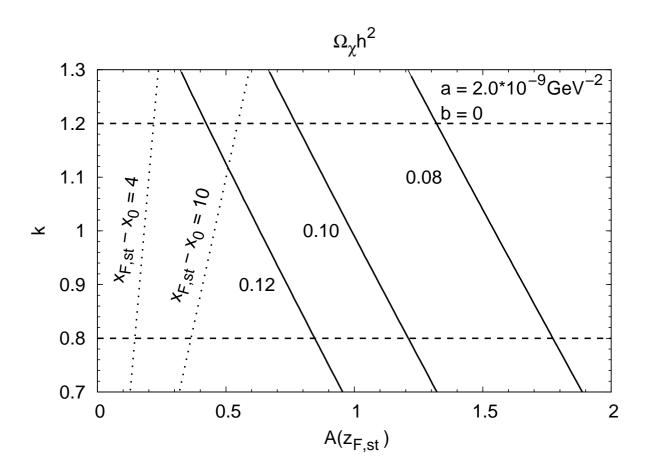
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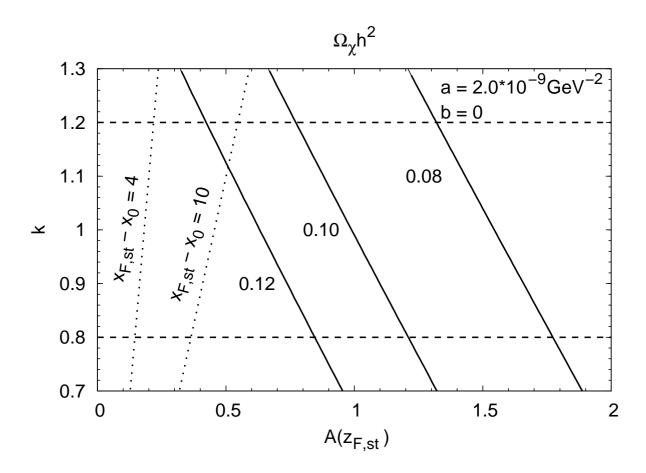


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Relative constraint on  $A(z_{F,st})$  weaker than that on  $\Omega_{\chi}h^2$ .  $H \gg H_{st}$  can be excluded from WIMP non–observation (Schelke et al.)

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- Detection needs ultrapure materials in deep–underground location; way to distinguish recoils from  $\beta, \gamma$  events; neutron screening; ...
- Is being pursued vigorously around the world!

#### **Direct WIMP detection: theory**

Counting rate given by

$$\frac{dR}{dQ} = AF^2(Q) \int_{v_{\min}}^{v_{\text{esc}}} \frac{f_1(v)}{v} dv$$

Q: recoil energy

 $A = \rho \sigma_0 / (2m_\chi m_r) = \text{const.: encodes particle physics}$ 

F(Q): nuclear form factor

v: WIMP velocity in lab frame

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In principle, can invert this relation to measure  $f_1(v)$ !

MD & C.L. Shan, astro-ph/0703651

$$f_1(v) = \mathcal{N} \left\{ -2Q \frac{d}{dQ} \left[ \frac{1}{F^2(Q)} \frac{dR}{dQ} \right] \right\}_{Q=2m_r^2 v^2/m_N}$$

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dR/dQ is approximately exponential: better work with logarithmic slope

### **Determining the logarithmic slope of** dR/dQ

 Good local observable: Average energy transfer  $\langle Q \rangle_i$  in *i*-th bin

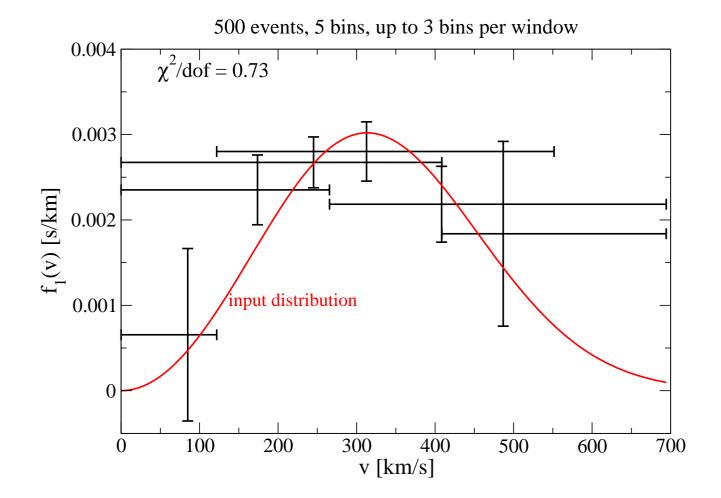
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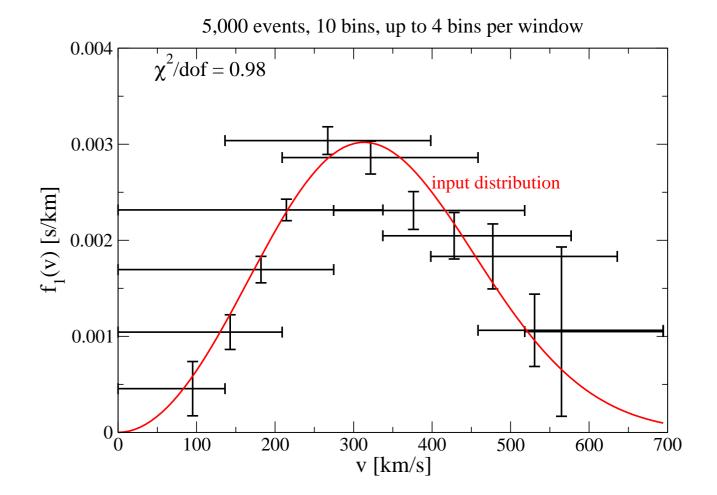
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- To maximize information: use overlapping bins ("windows")

#### **Recoil spectrum: prediction and simulated measurement**

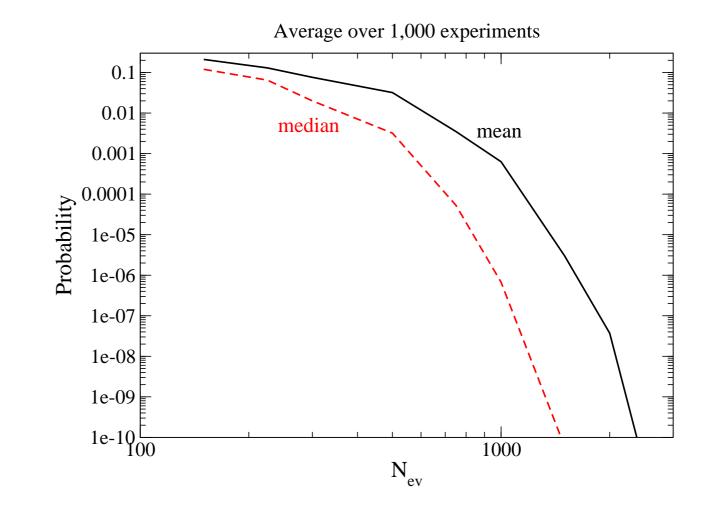


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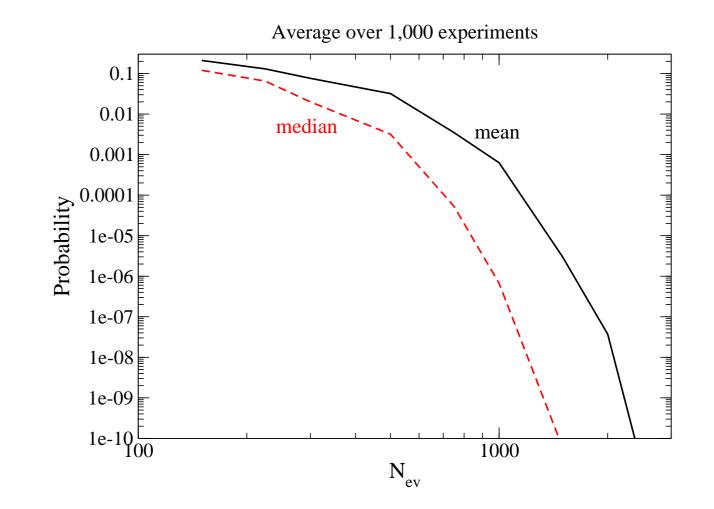


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#### **Statistical exclusion of constant** $f_1$



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\_Need several hundred events to begin direct reconstruction!\_

 $\langle v^n \rangle \equiv \int_0^\infty v^n f_1(v) dv$ 

$$\langle v^n \rangle \equiv \int_0^\infty v^n f_1(v) dv \\ \propto \int_0^\infty Q^{(n-1)/2} \frac{1}{F^2(Q)} \frac{dR}{dQ} dQ$$

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Can incorporate finite energy (hence velocity) threshold Moments are strongly correlated!

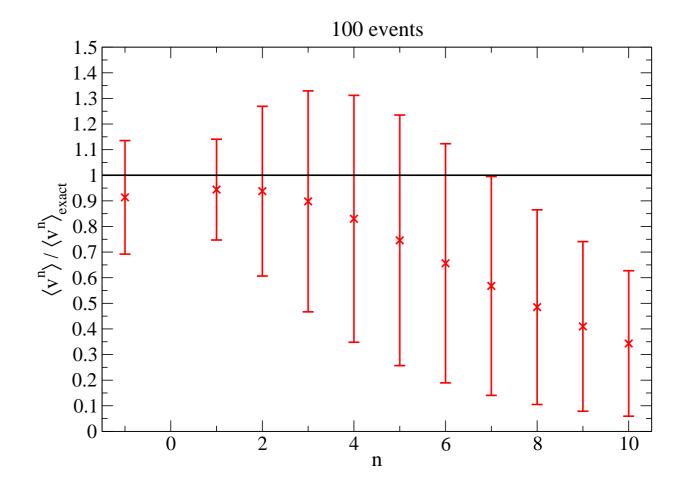
$$\begin{aligned} \langle v^n \rangle &\equiv \int_0^\infty v^n f_1(v) dv \\ &\propto \int_0^\infty Q^{(n-1)/2} \frac{1}{F^2(Q)} \frac{dR}{dQ} dQ \\ &\to \sum_{\text{events } a} \frac{Q_a^{(n-1)/2}}{F^2(Q_a)} \end{aligned}$$

Can incorporate finite energy (hence velocity) threshold

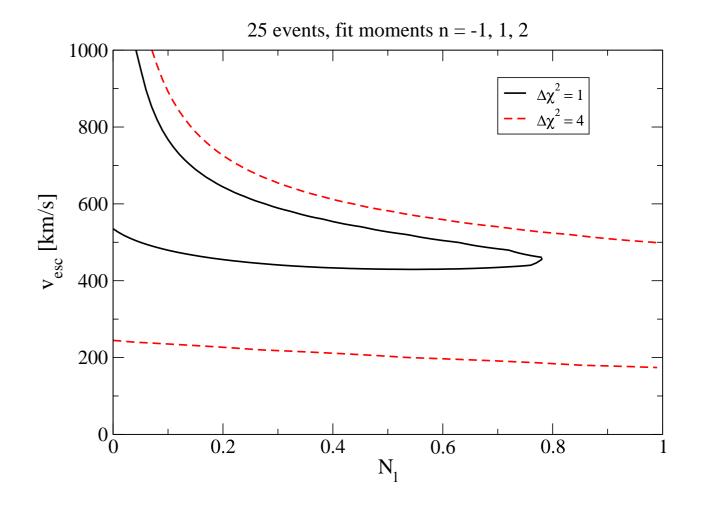
Moments are strongly correlated!

High moments, and their errors, are underestimated in "typical" experiment: get large contribution from large *Q* 

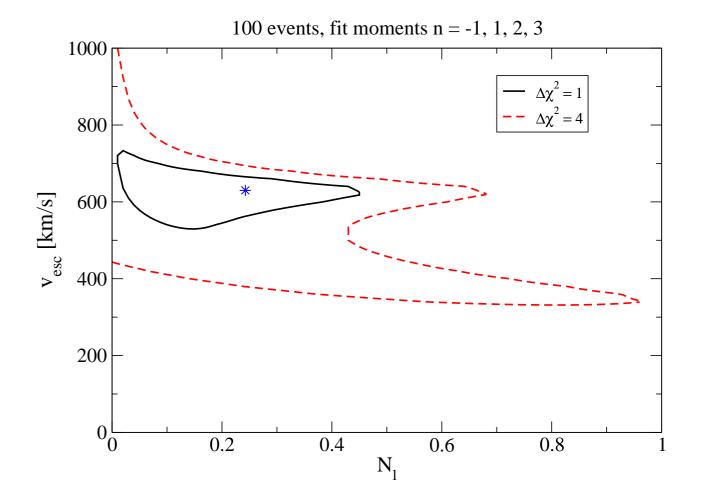
### **Determination of first 10 moments**



## **Constraining a "late infall" component**



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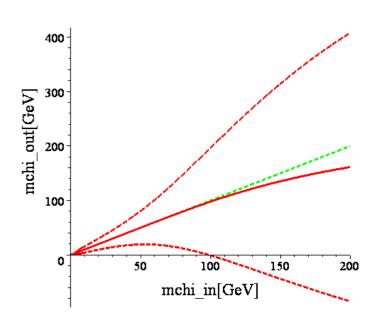
Learning from WIMPs - p. 26/32

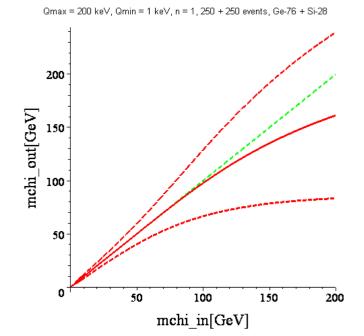
## **Determining the WIMP mass**

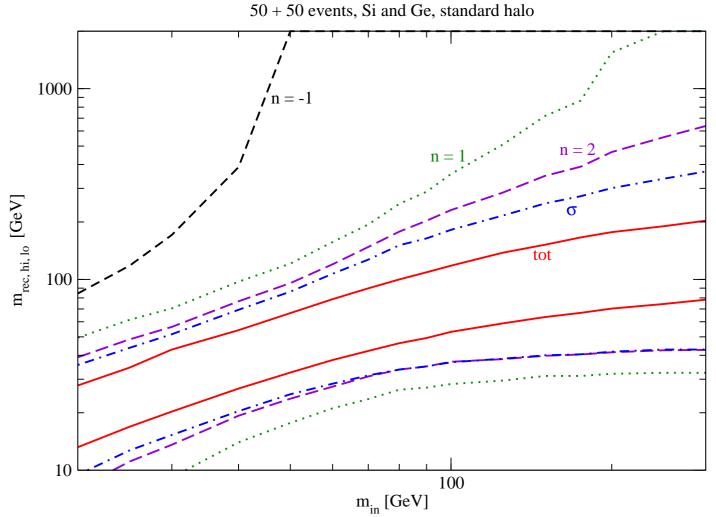
MD & C.L. Shan, in progress

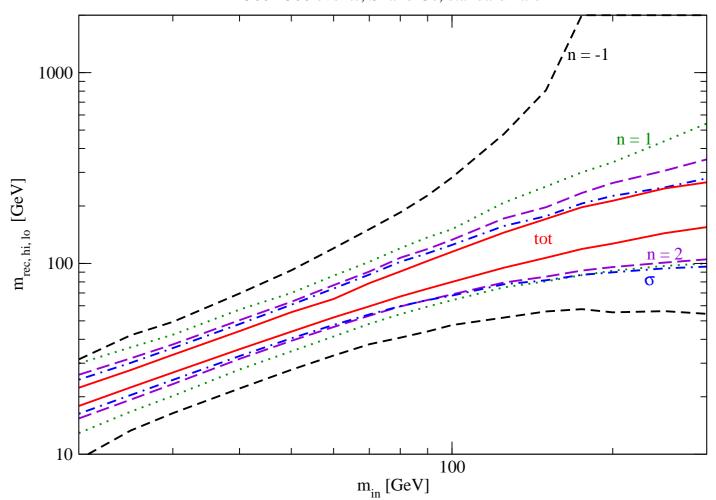
# Can determine $m_{\chi}$ from requirement that different targets yield same moments of $f_1$

Qmax = 200 keV, Qmin = 1 keV, n = 1, 25 + 25 events, Ge-76 + Si-28

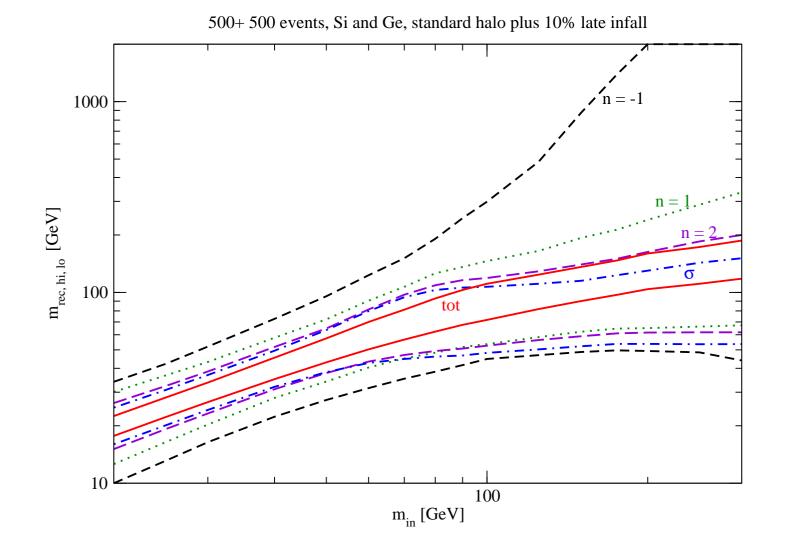


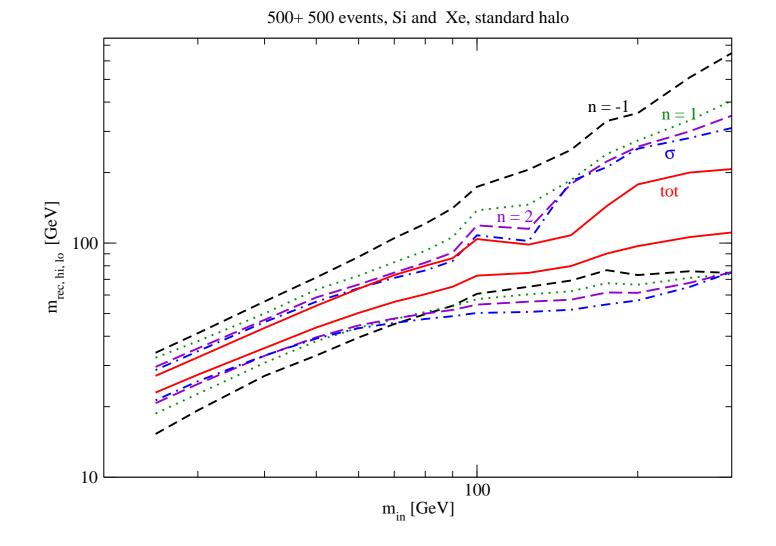






500+ 500 events, Si and Ge, standard halo





Learning from WIMPs – p. 31/32



Learning about the Early Universe:

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- Learning about WIMPs: Can determine  $m_{\chi}$  from moments of  $f_1$  measured with two different targets.