

Halo biasing with time renormalization group

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$$\delta(\mathbf{x}, t) = \frac{\rho(\mathbf{x}, t) - \bar{\rho}(t)}{\bar{\rho}(t)} \quad \langle \delta(\mathbf{k}) \delta(\mathbf{k}') \rangle \equiv \delta_D(\mathbf{k} + \mathbf{k}') P(k)$$

The **power spectrum** directly measures the fractional density contributions on different scales.

Observations	→	galaxy distribution
		\neq
Theory	→	mass distribution

The clustering of galaxies is different from that of the underlying matter distribution: $P_g(k) \neq P_m(k) \rightsquigarrow$ **bias**.

Since the galaxies form inside dark matter haloes, as a starting point we construct a model for dark haloes power spectrum P_h .

The large-scale structure observed today in the universe is believed to be the result of gravitational amplification of primordial fluctuations, caused by the interaction among cold dark matter particles.

Continuity equation + Euler equation + Poisson equation

$$\frac{\partial \delta_m}{\partial \tau} + \nabla \cdot [(1 + \delta_m)\mathbf{v}] = 0$$

$$\frac{\partial \mathbf{v}}{\partial \tau} + \mathcal{H}\mathbf{v} + (\mathbf{v} \cdot \nabla)\mathbf{v} = -\nabla\phi$$

$$\nabla^2\phi = \frac{3}{2} \mathcal{H}^2 \Omega_m \delta_m$$

Non-linear equations!

The distribution of dark matter haloes.

$$\frac{\partial \delta_h}{\partial \tau} + \nabla \cdot [(1 + \delta_h) \mathbf{v}_h] = 0$$

How?

1)

A **proto-halo** is the Lagrangian region of space that will collapse to form a halo at a certain redshift.

Unlike real haloes that undergo merging, by construction proto-haloes always preserve their identity.

2)

On scales much larger than the characteristic size of (and separation between) the proto-haloes, the density fluctuations traced by the centre-of-mass trajectories can be described with a continuous overdensity field δ_h .

But proto-haloes are discrete objects! \rightarrow shot noise ($\propto 1/N$)

$$\langle \delta_m(\mathbf{k}) \delta_h(\mathbf{k}') \rangle \equiv \delta_D(\mathbf{k} + \mathbf{k}') P_{mh}(k)$$

Time renormalization group (TRG) (Pietroni 2008)

Applying the “equations of motion” straightforwardly to the computation of the power spectrum, a hierarchy of equations is generated (like BBGKY), including higher orders (bispectrum $\langle \delta \delta \delta \rangle \dots$).

We close this system with a truncation at a definite order (**but** corrections will always be subdominant).

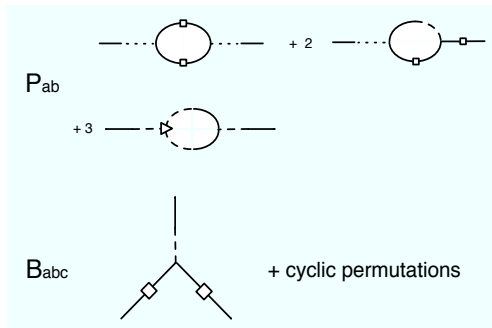
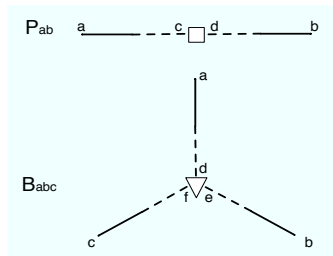
We have a system of differential equations which allows us to obtain the power spectrum at any time.

$$P(k)(\text{at early times} \Leftrightarrow \text{linear}) \mapsto \text{TRG} \mapsto P(k)(\text{today} \Leftrightarrow \text{non-linear})$$

In particle physics the RG equation describes the running of the coupling at different energy scales.



Time RG describes the running of the power spectrum at different time scales.



A non-local model for the halo bias at the initial conditions (adapted from Matsubara 1999)

The haloes are supposed to form from density peaks.

$$\delta_h(k) = (b_1 + b_2 k^2) \delta_m(k), \quad \mathbf{v}_h(k) = \mathbf{v}(k)$$

$$P_{mh}(k) = (b_1 + b_2 k^2) P_m(k) e^{-k^2 R^2/2}$$

- ◇ the exponential function accounts for the exclusion effects of the haloes, due to their finite size;
- ◇ P_m is from linear theory;
- ◇ we can also compute the **bispectrum** for the haloes, which is important because the haloes density field is not gaussian!

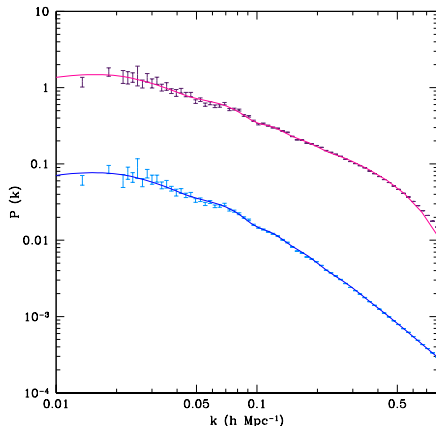
$$P_{mh}(k) = (b_1 + b_2 k^2) P_m(k) e^{-k^2 R^2/2}$$

We fit the parameters of the model from P_{mh} measured at a redshift $z=50$ from the N-body simulations by Pillepich et al. 2010, for haloes with a mass $M > 10^{13} M_\odot$:

$$\# \quad b_1 = 19.2 \pm 0.3$$

$$\# \quad b_2 = 467 \pm 66 \text{ Mpc}^2 h^{-2}$$

$$\# \quad R = 2.6 \pm 0.4 \text{ Mpc}/h$$



AIM: a power spectrum which accounts for **non-linear evolution** and **halo biasing**

DM particles equations

+

DM haloes equation

Bias model (IC)



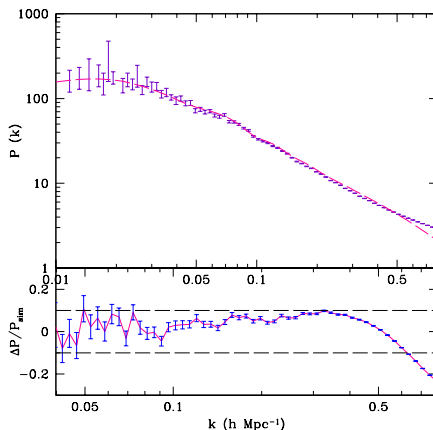
TRG



Power spectrum P_{mh}

Power spectrum today ($z = 0$)

TRG result versus N-body simulations (box size: $1200 h^{-1}$ Mpc)



Difference within 10% : $k \sim 0.6 h/\text{Mpc}$

Conclusions

- The TRG method allows us to obtain a cross PS with good precision up to non-linear scales.
- The adopted model for the initial bias has proven successful.
- An analytic formalism is also available.
- Future work: redshift space distortion.

TRG equations

$$\begin{aligned}
 \partial_\eta P_{ab}(k) &= -\Omega_{ac}P_{cb}(k) - \Omega_{bc}P_{ac}(k) \\
 &\quad + e^\eta \frac{4\pi}{k} [I_{acd,bcd}(k) + I_{bcd,acd}(k)] , \\
 \partial_\eta I_{acd,bef}(k) &= -\Omega_{bg}I_{acd,gef}(k) - \Omega_{eg}I_{acd,bgf}(k) \\
 &\quad - \Omega_{fg}I_{acd,beg}(k) + 2e^\eta A_{acd,bef}(k) .
 \end{aligned}$$

$$\begin{aligned}
 I_{acd,bef}(k) &\equiv \int_{k/2}^{\infty} dq \, q \int_{|q-k|}^q dp \, p \, \frac{1}{2} \left[\tilde{\gamma}_{acd}(k, q, p) \tilde{B}_{bef}(k, q, p) + (q \leftrightarrow p) \right] . \\
 A_{acd,bef}(k) &\equiv \int_{k/2}^{\infty} dq \, q \int_{|q-k|}^q dp \, p \, \frac{1}{2} \{ \tilde{\gamma}_{acd}(k, q, p) [\tilde{\gamma}_{bgh}(k, q, p) P_{ge}(q) P_{hf}(p) \\
 &\quad + \tilde{\gamma}_{egh}(q, p, k) P_{gf}(p) P_{hb}(k) + \tilde{\gamma}_{fgh}(p, k, q) P_{gb}(k) P_{he}(q)] + (q \leftrightarrow p) \} .
 \end{aligned}$$

Bispectrum expressions

$$B_{hhh}(k_1, k_2, k_3) = (b_1 + b_2 k_1^2)(b_1 + b_2 k_2^2)(b_1 + b_2 k_3^2) B_m(k_1, k_2, k_3),$$

$$B_{mhh}(k_1, k_2, k_3) = (b_1 + b_2 k_2^2)(b_1 + b_2 k_3^2) B_m(k_1, k_2, k_3),$$

$$B_{mmh}(k_1, k_2, k_3) = (b_1 + b_2 k_3^2) B_m(k_1, k_2, k_3),$$

The matter bispectrum at tree level is:

$$B_m(k_1, k_2, k_3) = 2 \left[\frac{1}{2} + \frac{1}{2} \left(\frac{k_1}{k_2} + \frac{k_2}{k_1} \right) \frac{\vec{k}_1 \cdot \vec{k}_2}{k_1 k_2} + \frac{1}{2} \left(\frac{\vec{k}_1 \cdot \vec{k}_2}{k_1 k_2} \right)^2 \right] P(k_1, z_{in}) P(k_2, z_{in}) + \text{cycl.}$$

Shot-noise problem

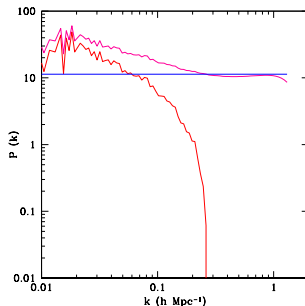
Selected mass range: $M > 1.24 \cdot 10^{13} M_{\odot}$
(haloes with more than 100 particles)



Number of haloes: 613343



Shot noise: $2817.35 \text{ (Mpc}/h)^3$



N-body simulation (Pillepich, Porciani, Hahn 2009)

Specifics:

- # number of particles: 1024^3 ;
- # box size: $1200 h^{-1}$ Mpc;
- # resolution: 512;
- # initial redshift: $z = 50$;
- # cosmology: $h = 0.701, \sigma_8 = 0.817, n_s = 0.96,$
 $\Omega_m = 0.279, \Omega_b = 0.0462, \Omega_\Lambda = 0.721.$