

Learning from WIMPs

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Introduction: WIMPs as Dark Matter

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- Models of structure formation, X ray temperature of clusters of galaxies, ...
- **Cosmic Microwave Background anisotropies (WMAP)**
imply $\Omega_{\text{DM}}h^2 = 0.105^{+0.007}_{-0.013}$ Spergel et al., astro-ph/0603449

Weakly Interacting Massive Particles (WIMPs)

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- In standard cosmology, roughly weak cross section automatically gives roughly right relic density for thermal WIMPs! (On logarithmic scale)
- Roughly weak interactions may allow both *direct* and *indirect* detection of WIMPs

WIMP production

Let χ be a generic DM particle, n_χ its number density (unit: GeV^3). Assume $\chi = \bar{\chi}$, i.e. $\chi\chi \leftrightarrow \text{SM particles}$ is possible, but single production of χ is forbidden by some symmetry.

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Evolution of n_χ determined by **Boltzmann equation**:

$$\frac{dn_\chi}{dt} + 3Hn_\chi = -\langle\sigma_{\text{ann}}v\rangle (n_\chi^2 - n_{\chi,\text{eq}}^2)$$

$H = \dot{R}/R$: Hubble parameter

$\langle\dots\rangle$: Thermal averaging

$\sigma_{\text{ann}} = \sigma(\chi\chi \rightarrow \text{SM particles})$

v : relative velocity between χ 's in their cms

$n_{\chi,\text{eq}}$: χ density in full equilibrium

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Gives

$$\Omega_\chi h^2 \propto \frac{1}{\langle v \sigma_{\text{ann}} \rangle} \sim 0.1 \text{ for } \sigma_{\text{ann}} \sim \text{pb}$$

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Can we test these assumptions, if Ω_χ and “all” particle physics properties of χ are known?

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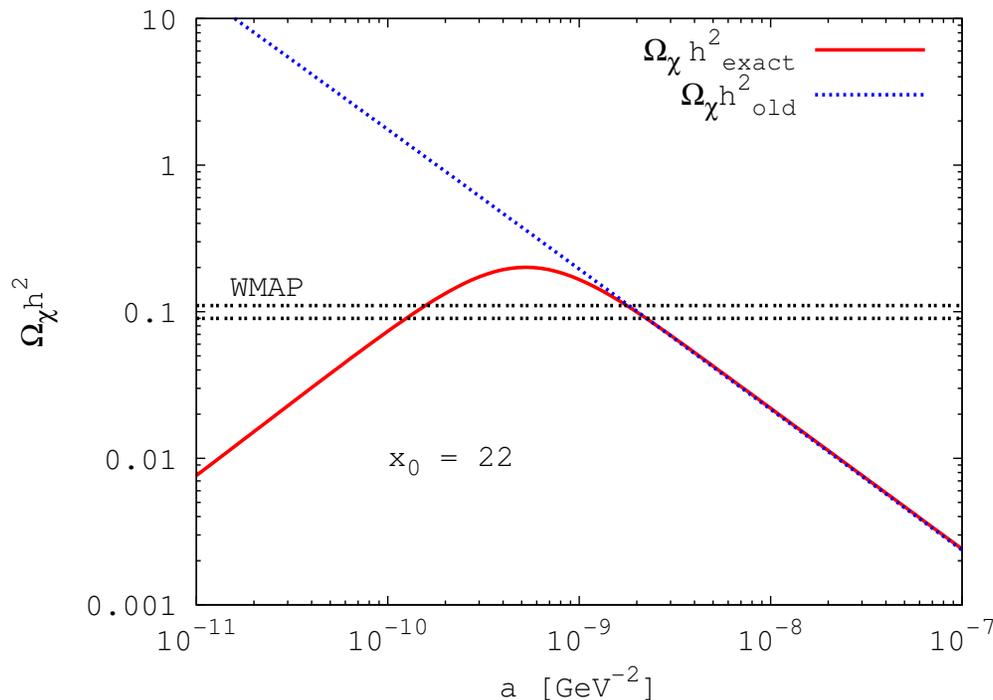
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Low temperature scenario (cont.'d)

Using explicit form of H , $Y_{\chi,\text{eq}}$, Boltzmann eq. becomes

$$\frac{dY_{\chi}}{dx} = -f \left(a + \frac{6b}{x} \right) x^{-2} \left(Y_{\chi}^2 - cx^3 e^{-2x} \right) .$$

$$f = 1.32 m_{\chi} M_{\text{Pl}} \sqrt{g_*}, \quad c = 0.0210 g_{\chi}^2 / g_*^2$$

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For $T_0 \ll T_F$: Annihilation term $\propto Y_{\chi}^2$ negligible: defines 0-th order solution $Y_0(x)$, with

$$Y_0(x \rightarrow \infty) = fc \left[\frac{a}{2} x_R e^{-2x_R} + \left(\frac{a}{4} + 3b \right) e^{-2x_R} \right] .$$

Note: $\Omega_{\chi} h^2 \propto \sigma_{\text{ann}}$ in this case!

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For intermediate temperatures, $T_0 \lesssim T_F$: Define 1st-order solution

$$Y_1 = Y_0 + \delta .$$

$\delta < 0$ describes pure annihilation:

$$\frac{d\delta}{dx} = -f \left(a + \frac{6b}{x} \right) \frac{Y_0(x)^2}{x^2} .$$

$\delta(x)$ can be calculated analytically: $\delta \propto \sigma_{\text{ann}}^3$

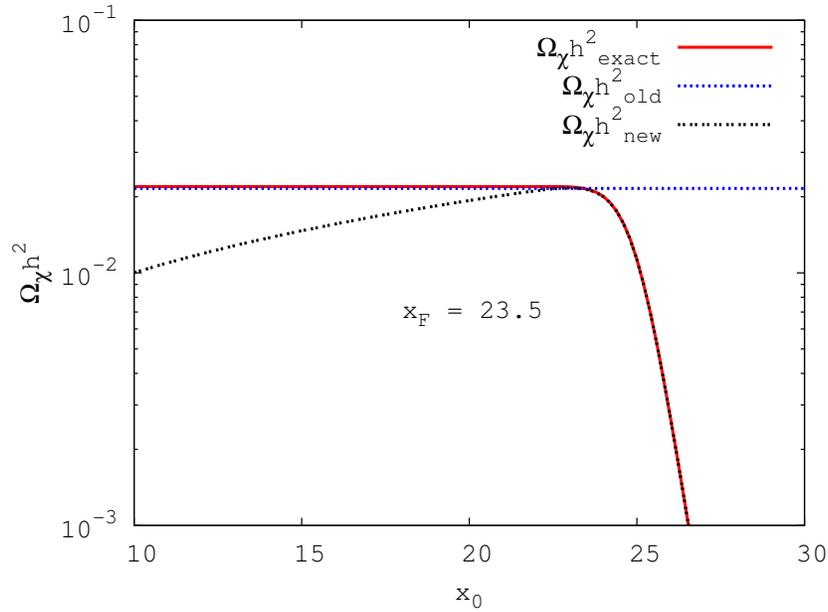
Low temperature scenario (cont.'d)

Get good results for $\Omega_\chi h^2$ for all $T_0 \leq T_F$ through “resummation”:

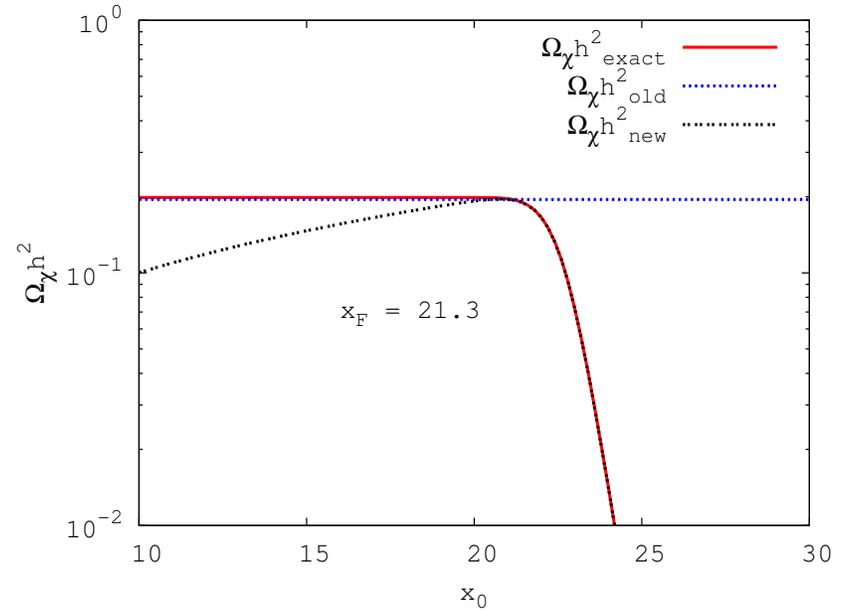
$$Y_1 = Y_0 \left(1 + \frac{\delta}{Y_0} \right) \simeq \frac{Y_0}{1 - \delta/Y_0} \equiv Y_{1,r}$$

$Y_{1,r} \propto 1/\sigma_{\text{ann}}$ for $|\delta| \gg Y_0$ MD, Imminniyaz, Kakizaki, hep-ph/0603165

Numerical comparison: $b = 0$

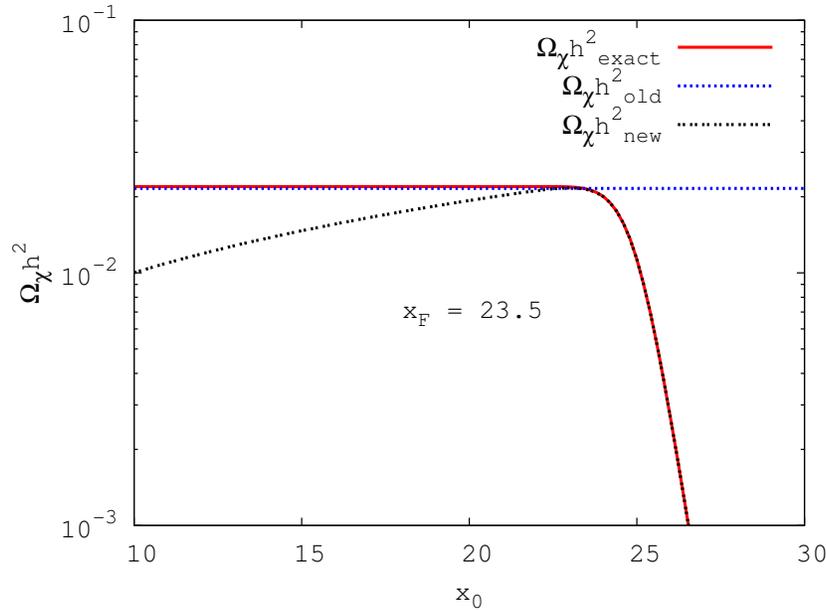


$$a = 10^{-8} \text{ GeV}^{-2}$$

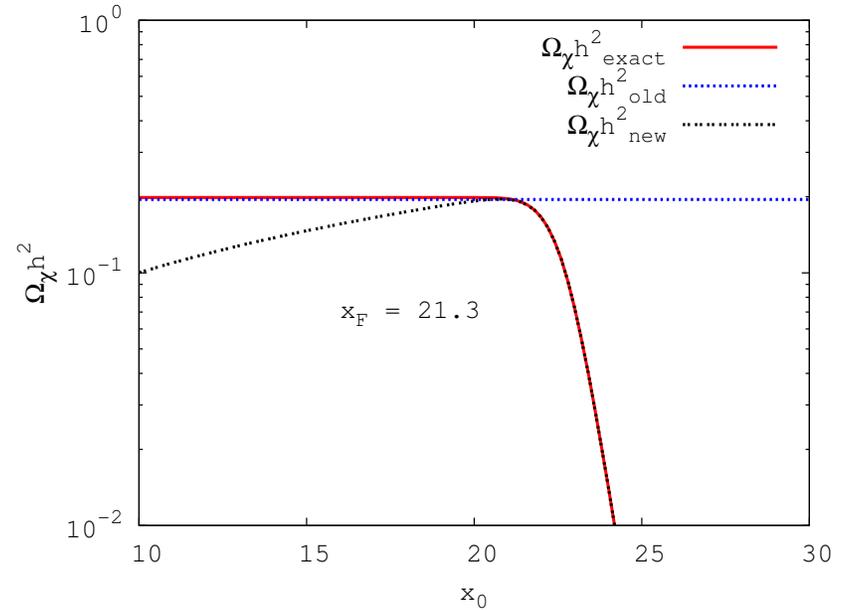


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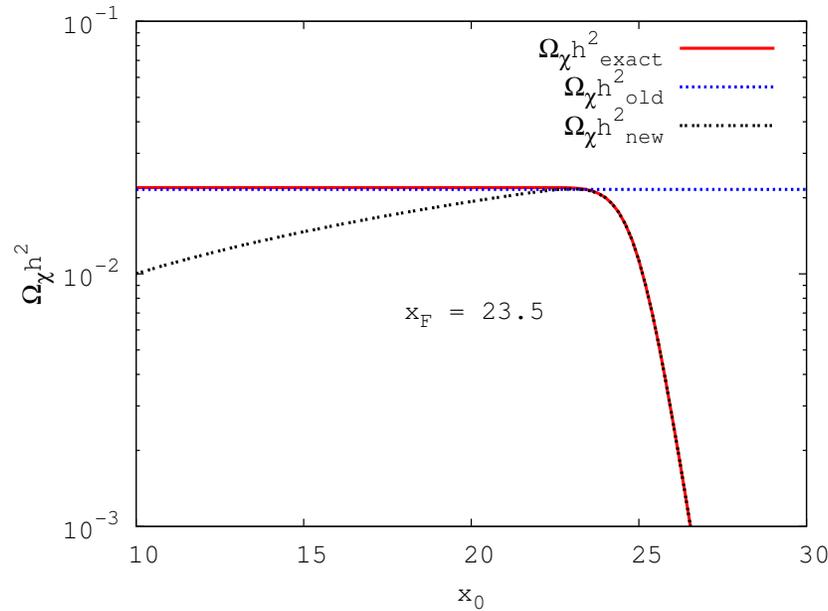
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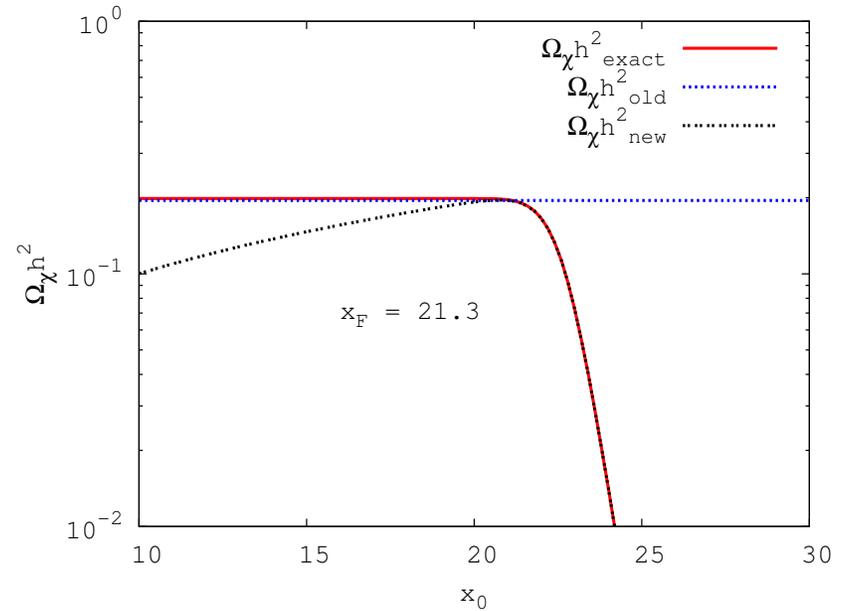
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Can extend validity of new solution to all T , including $T \gg T_0$, by using $\Omega_\chi(T_{\text{max}})$ if $T_0 > T_{\text{max}} \simeq T_F$

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Note: $\Omega_\chi(T_0) \leq \Omega_\chi(T_0 \gg T_F)$

Application: lower bound on T_0 for thermal WIMP

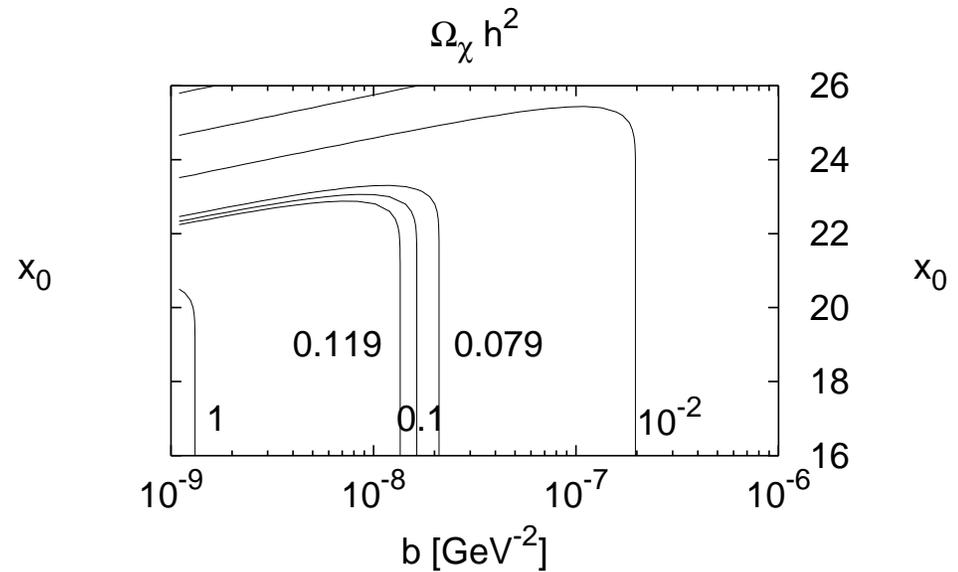
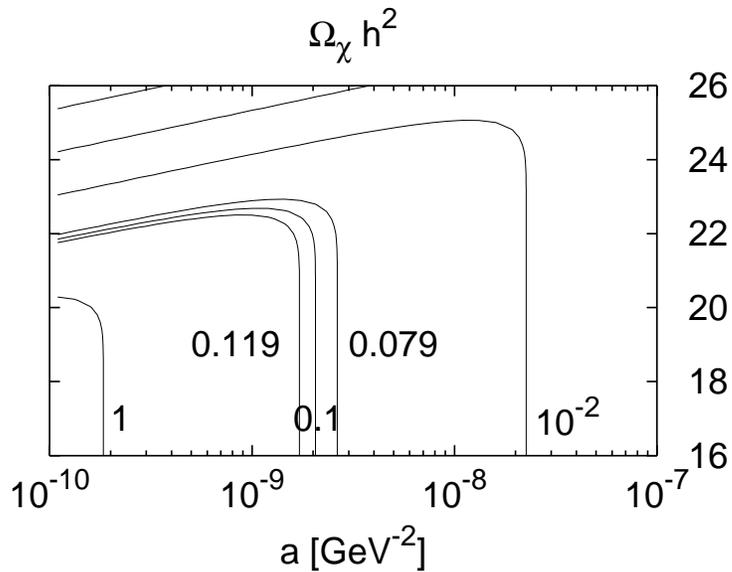
MD, Imminniyaz, Kakizaki, arXiv:0704.1590 [hep-ph]

If $n_\chi(T_0) = 0$, demanding $\Omega_\chi h^2 \simeq 0.1$ imposes lower bound on T_0 :

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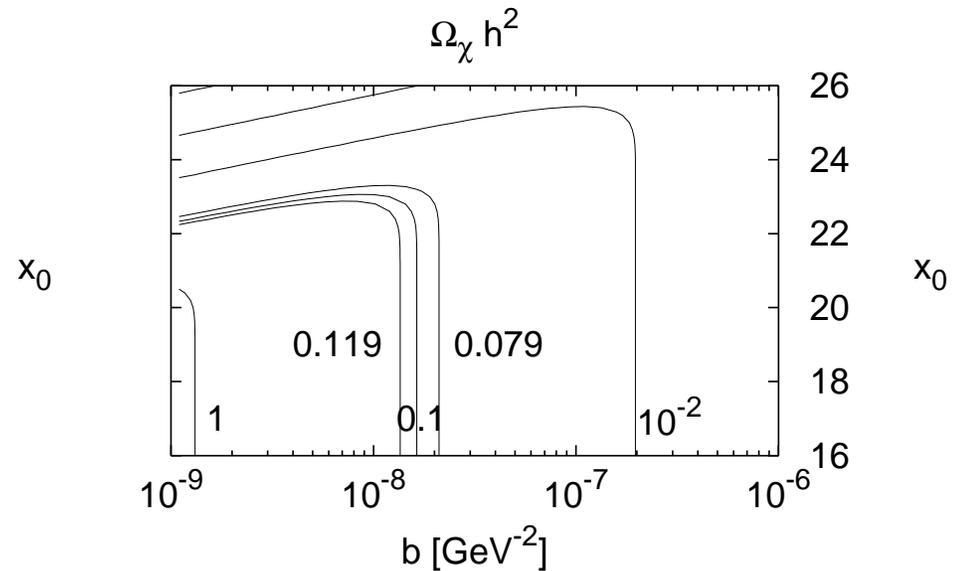
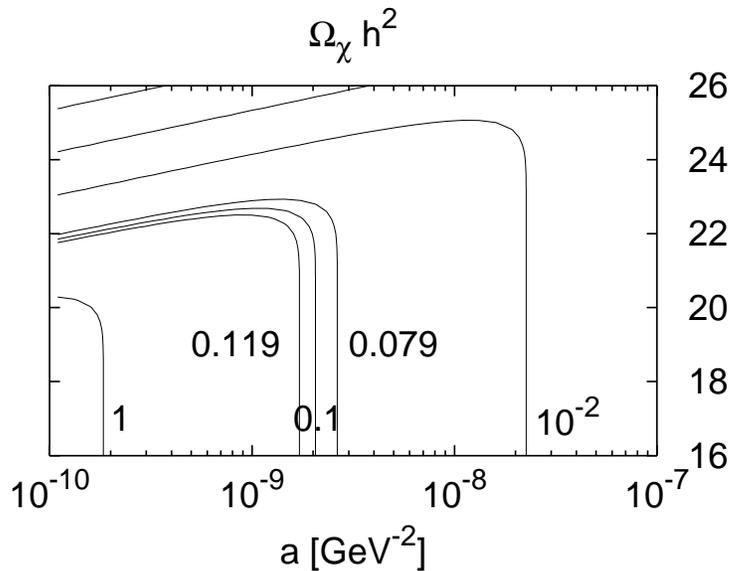
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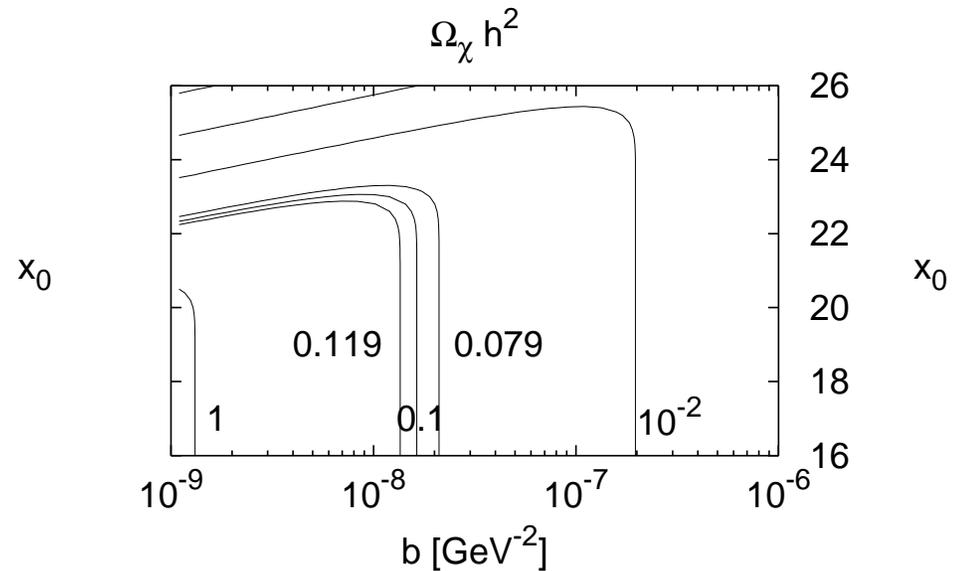
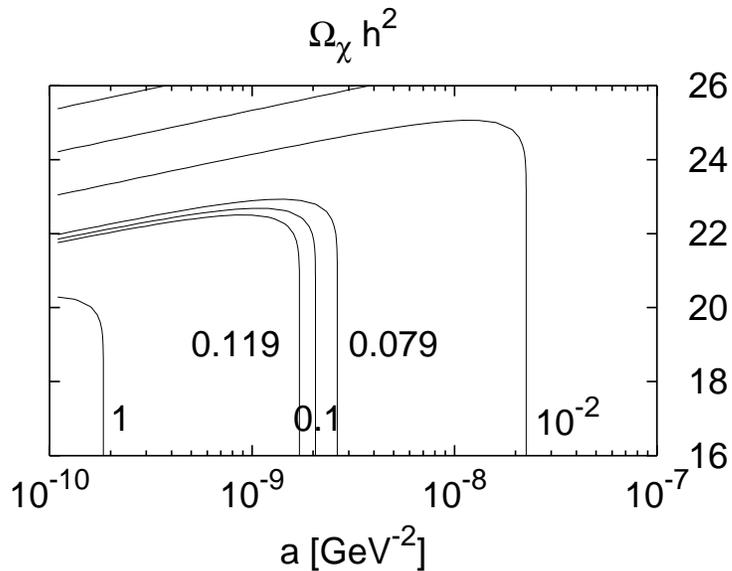


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If $T_0 \simeq m_\chi/22$: Get right $\Omega_\chi h^2$ for wide range of cross sections!

Constraining $H(T)$

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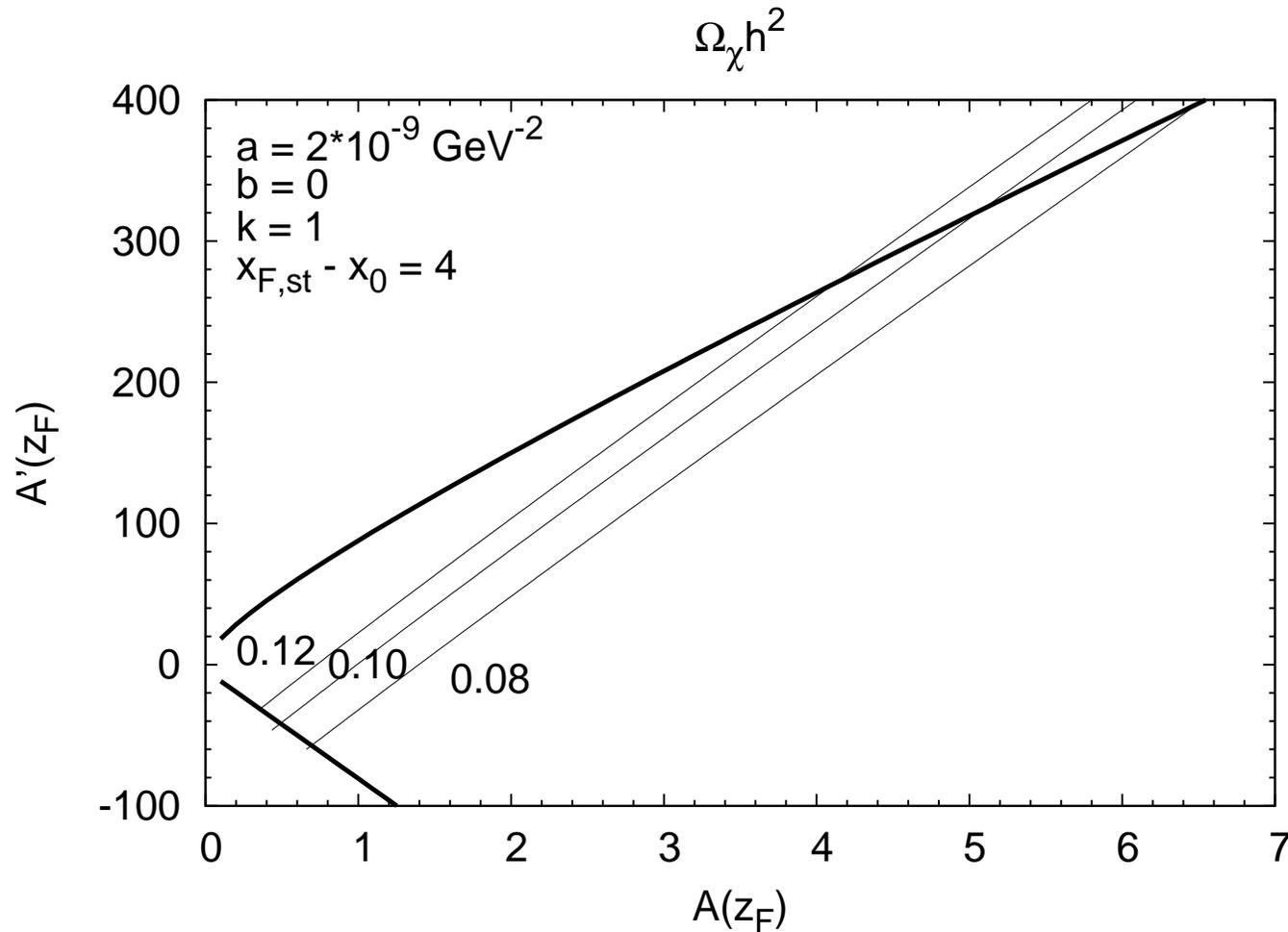
- Successful BBN $\implies k \equiv A(z \rightarrow 0) = 1.0 \pm 0.2$

Constraining $H(T)$ (cont.d)

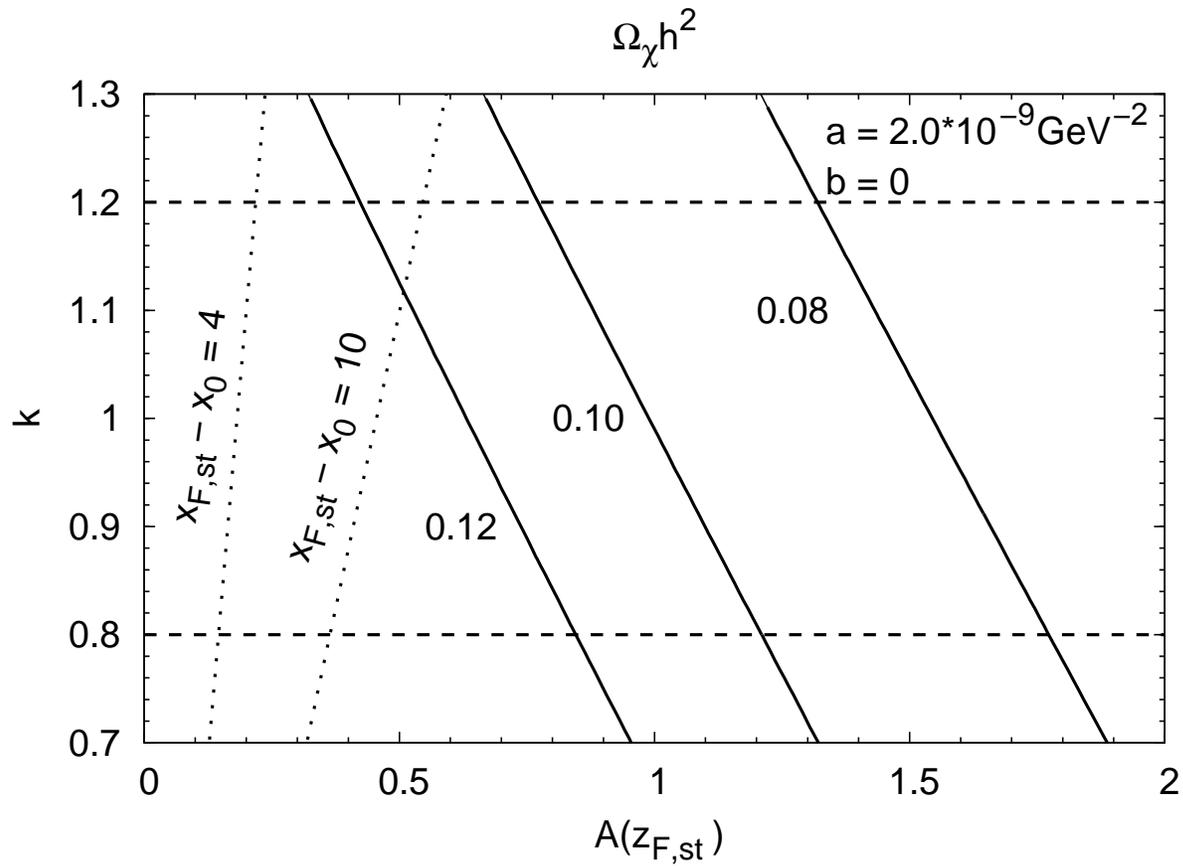
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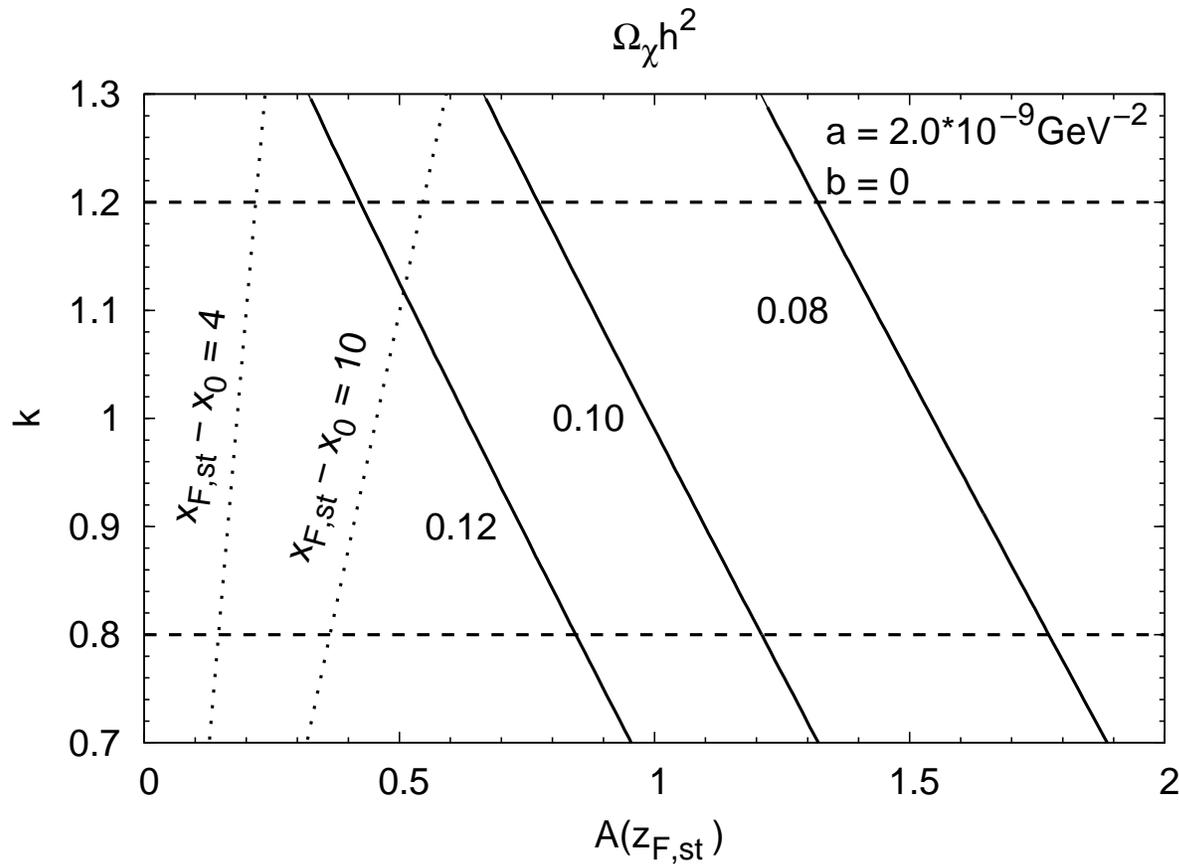
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Relative constraint on $A(z_{F,st})$ weaker than that on $\Omega_\chi h^2$.

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- Is being pursued vigorously around the world!

Direct WIMP detection: theory

Counting rate given by

$$\frac{dR}{dQ} = AF^2(Q) \int_{v_{\min}}^{v_{\text{esc}}} \frac{f_1(v)}{v} dv$$

Q : recoil energy

$A = \rho\sigma_0 / (2m_\chi m_r) = \text{const.}$: encodes particle physics

$F(Q)$: nuclear form factor

v : WIMP velocity in lab frame

$$v_{\min}^2 = m_N Q / (2m_r^2)$$

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In principle, can invert this relation to measure $f_1(v)$!

Direct reconstruction of f_1

MD & C.L. Shan, astro-ph/0703651

$$f_1(v) = \mathcal{N} \left\{ -2Q \frac{d}{dQ} \left[\frac{1}{F^2(Q)} \frac{dR}{dQ} \right] \right\}_{Q=2m_r^2 v^2 / m_N}$$

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dR/dQ is approximately exponential: better work with logarithmic slope

Determining the logarithmic slope of dR/dQ

- Good local observable: Average energy transfer $\langle Q \rangle_i$ in i -th bin

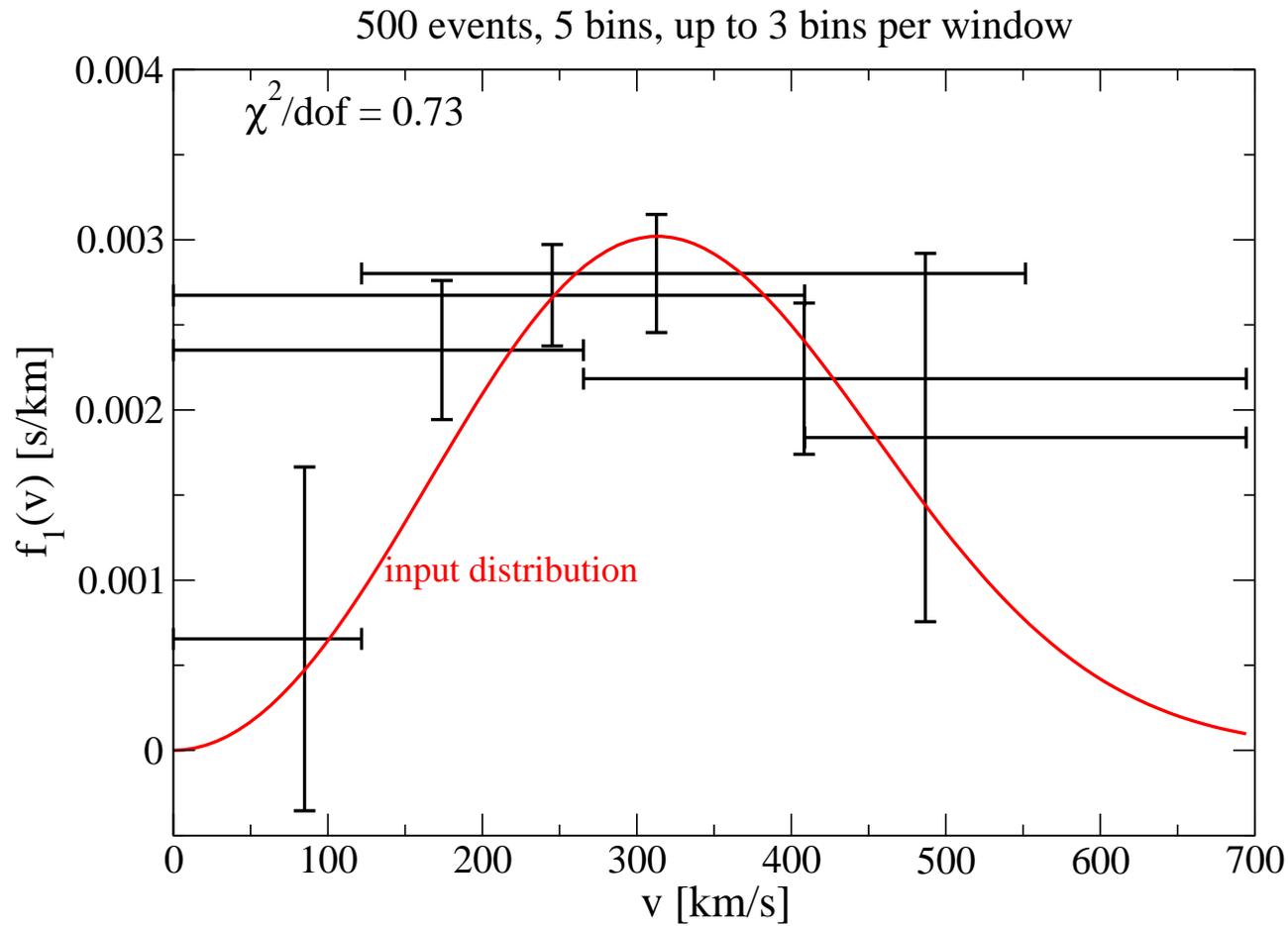
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- Good local observable: Average energy transfer $\langle Q \rangle_i$ in i -th bin
- Stat. error on slope $\propto (\text{bin width})^{-1.5} \implies$ **need large bins**

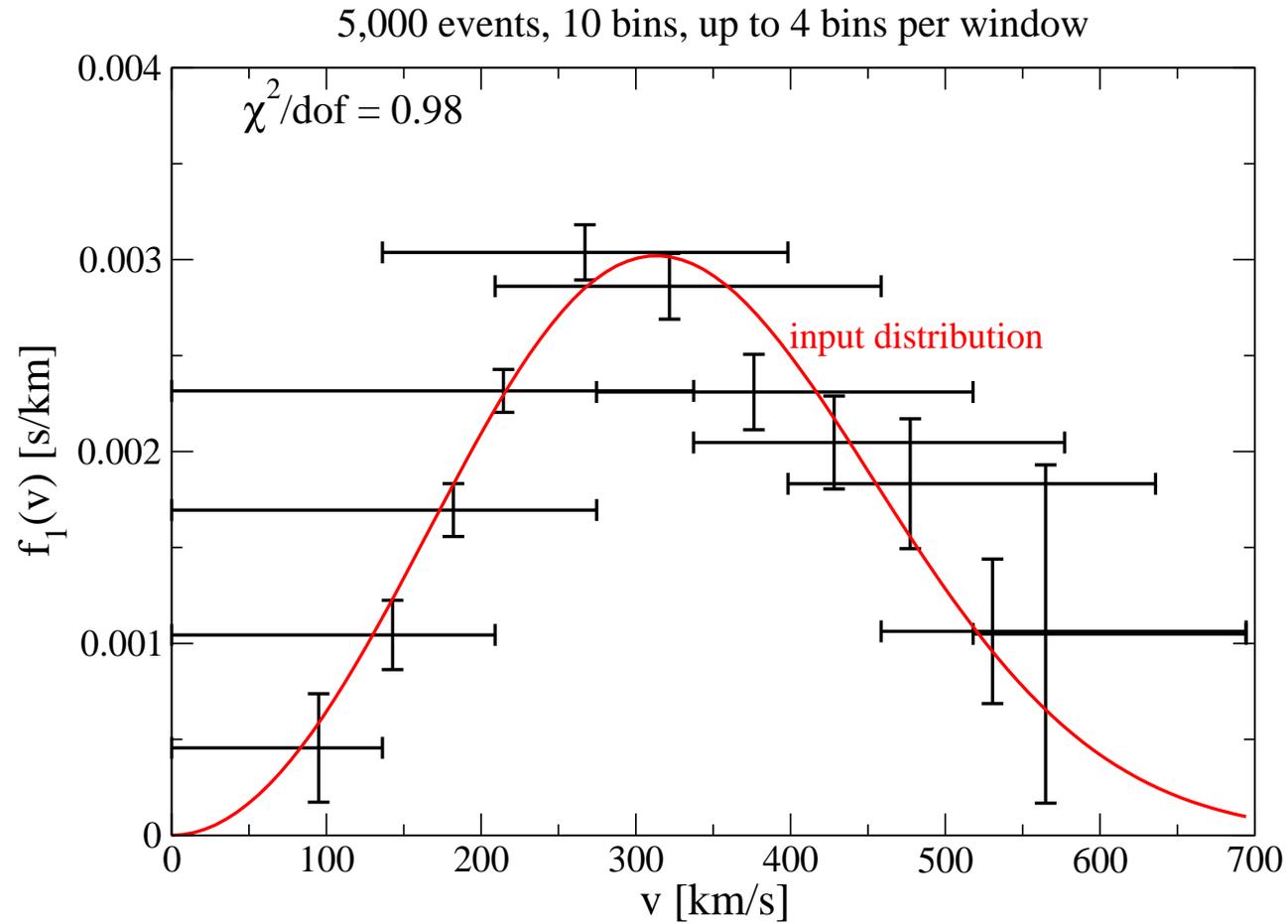
Determining the logarithmic slope of dR/dQ

- Good local observable: Average energy transfer $\langle Q \rangle_i$ in i -th bin
- Stat. error on slope $\propto (\text{bin width})^{-1.5} \implies$ **need large bins**
- To maximize information: **use overlapping bins** (“windows”)

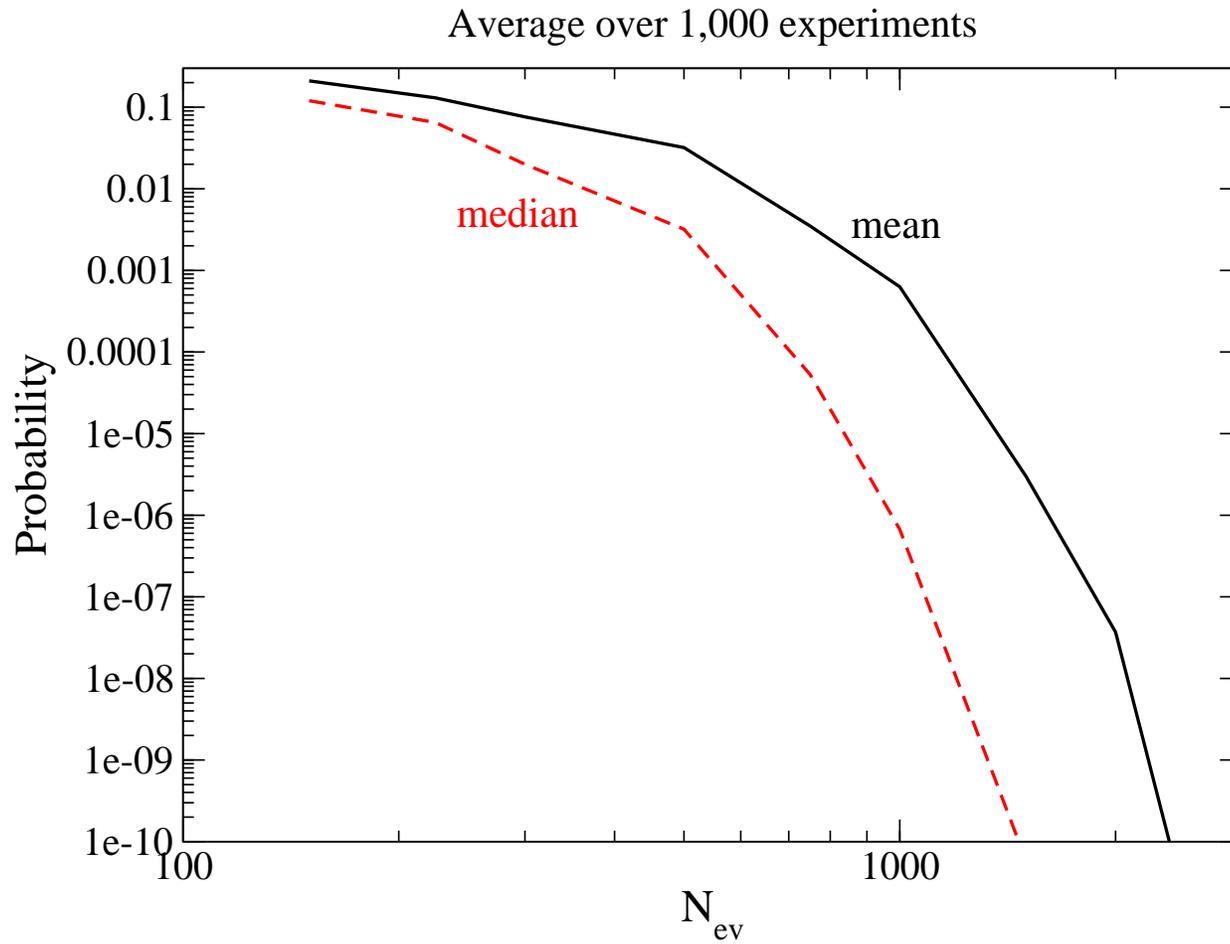
Recoil spectrum: prediction and simulated measurement



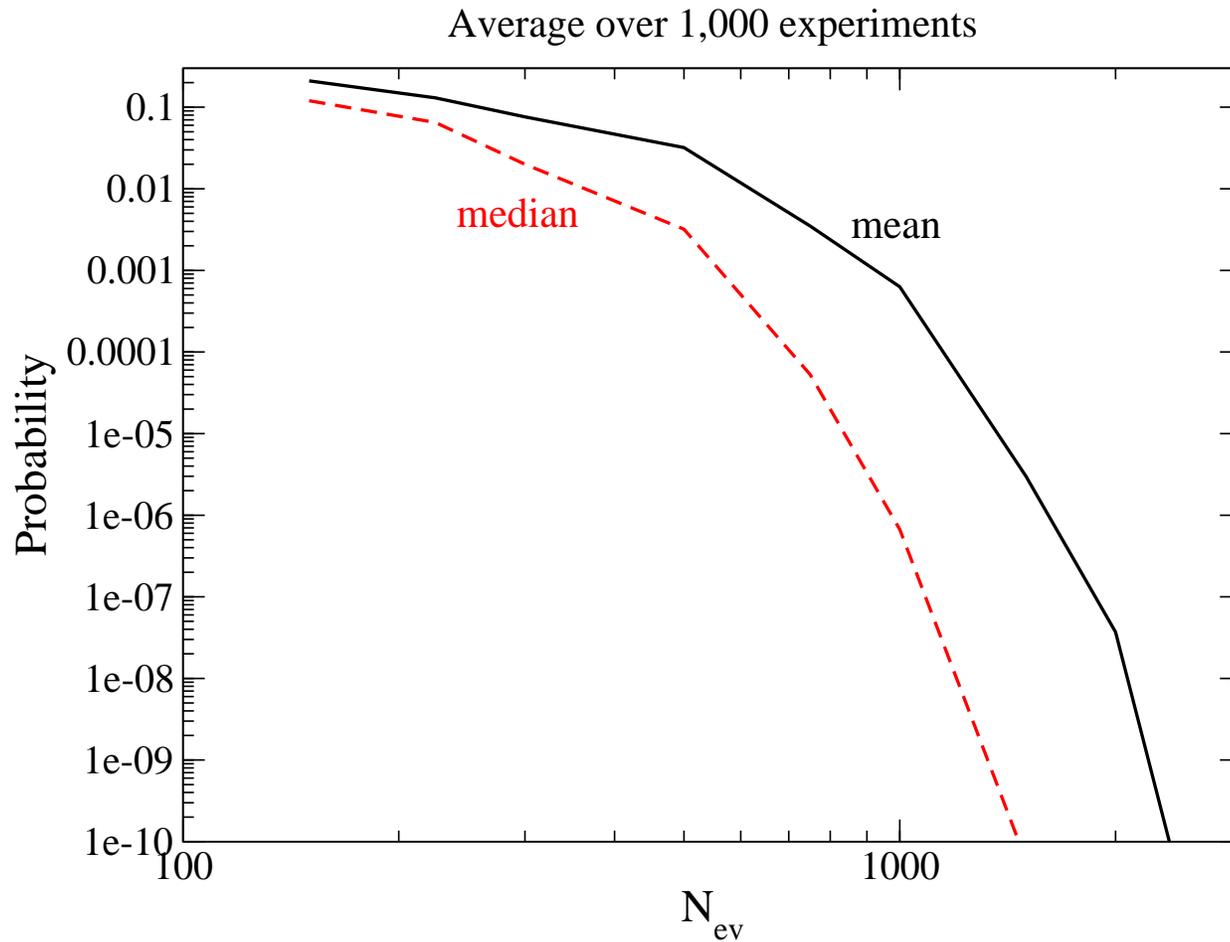
Recoil spectrum: prediction and simulated measurement



Statistical exclusion of constant f_1



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Need several hundred events to begin direct reconstruction!

Determining moments of f_1

$$\langle v^n \rangle \equiv \int_0^\infty v^n f_1(v) dv$$

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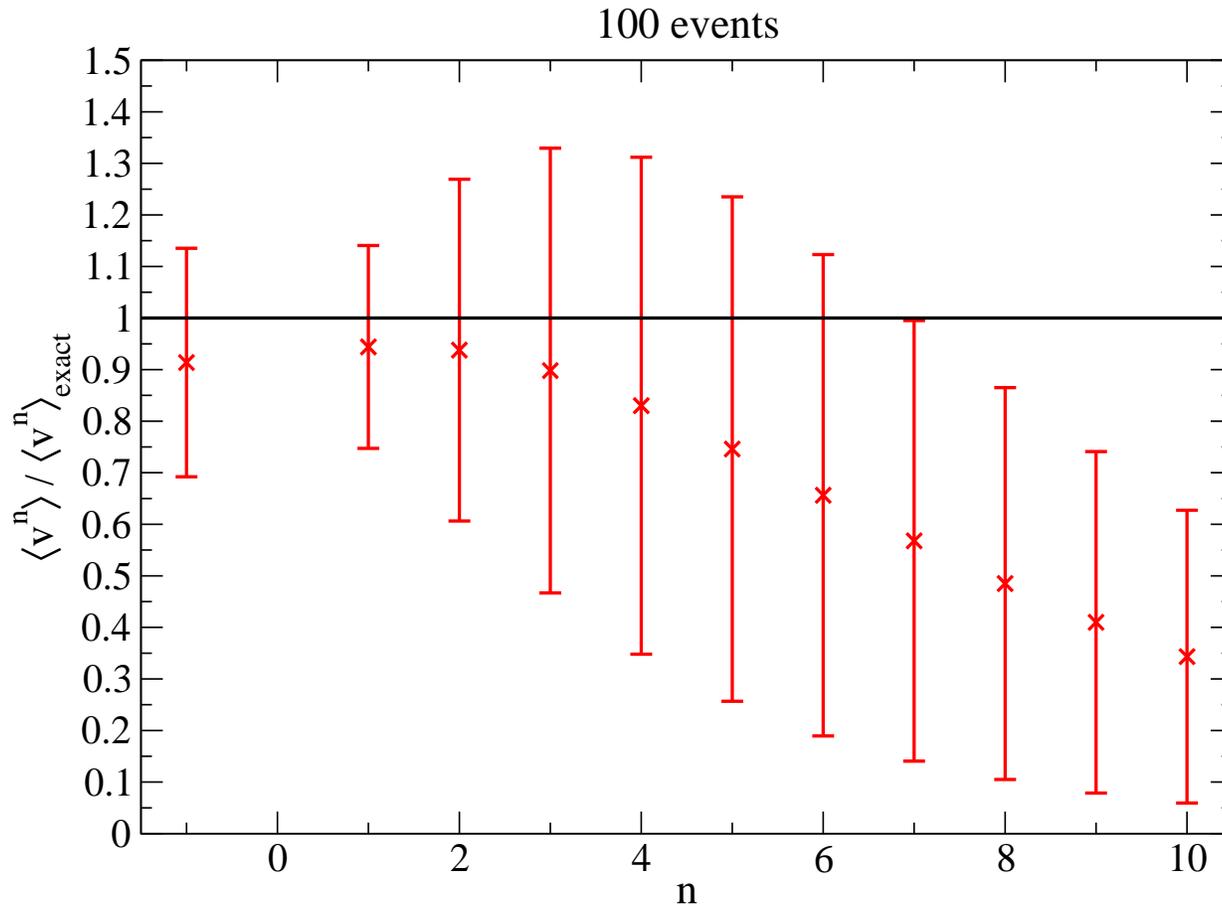
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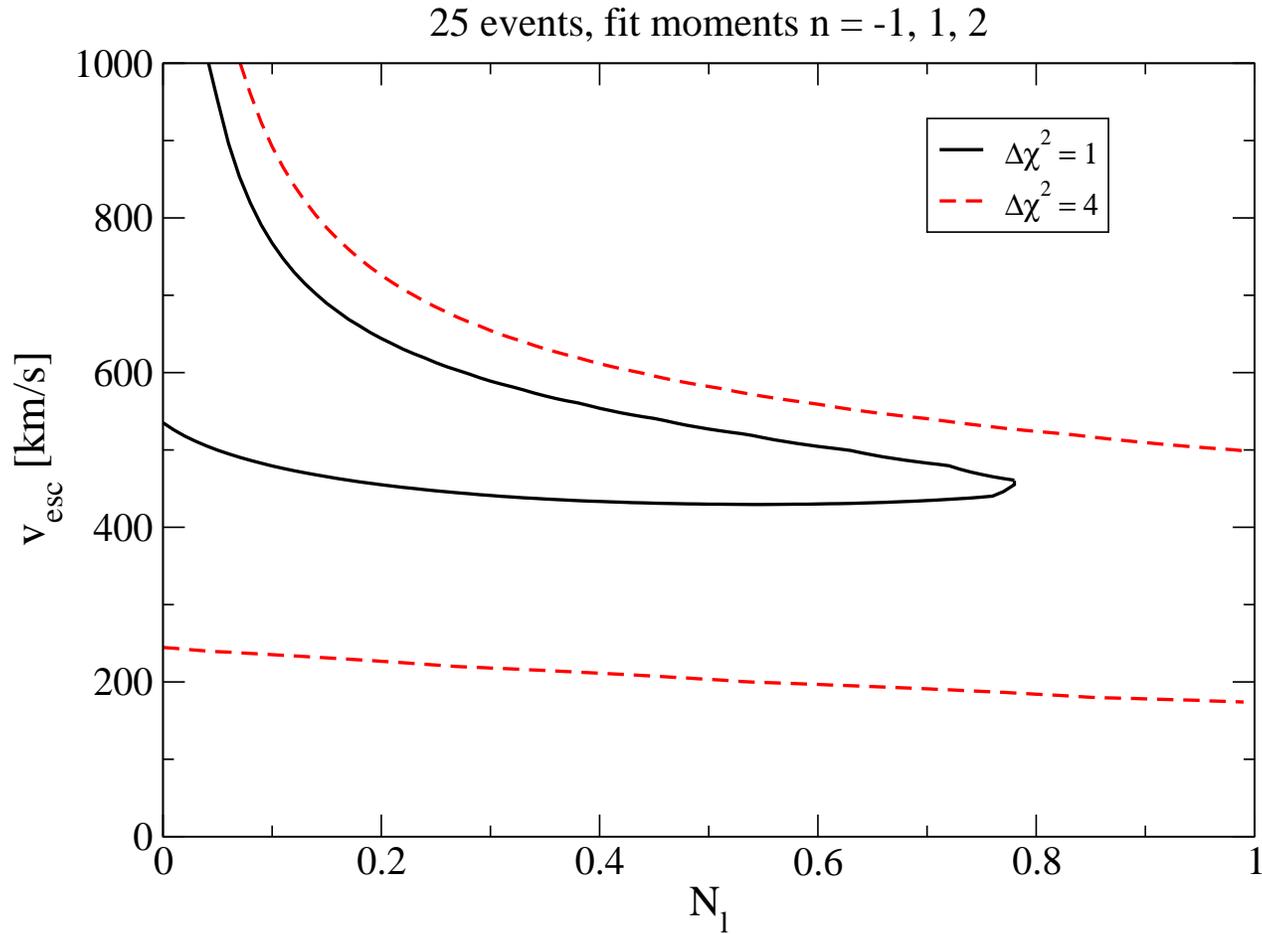
Moments are strongly correlated!

High moments, and their errors, are underestimated in “typical” experiment: get large contribution from large Q

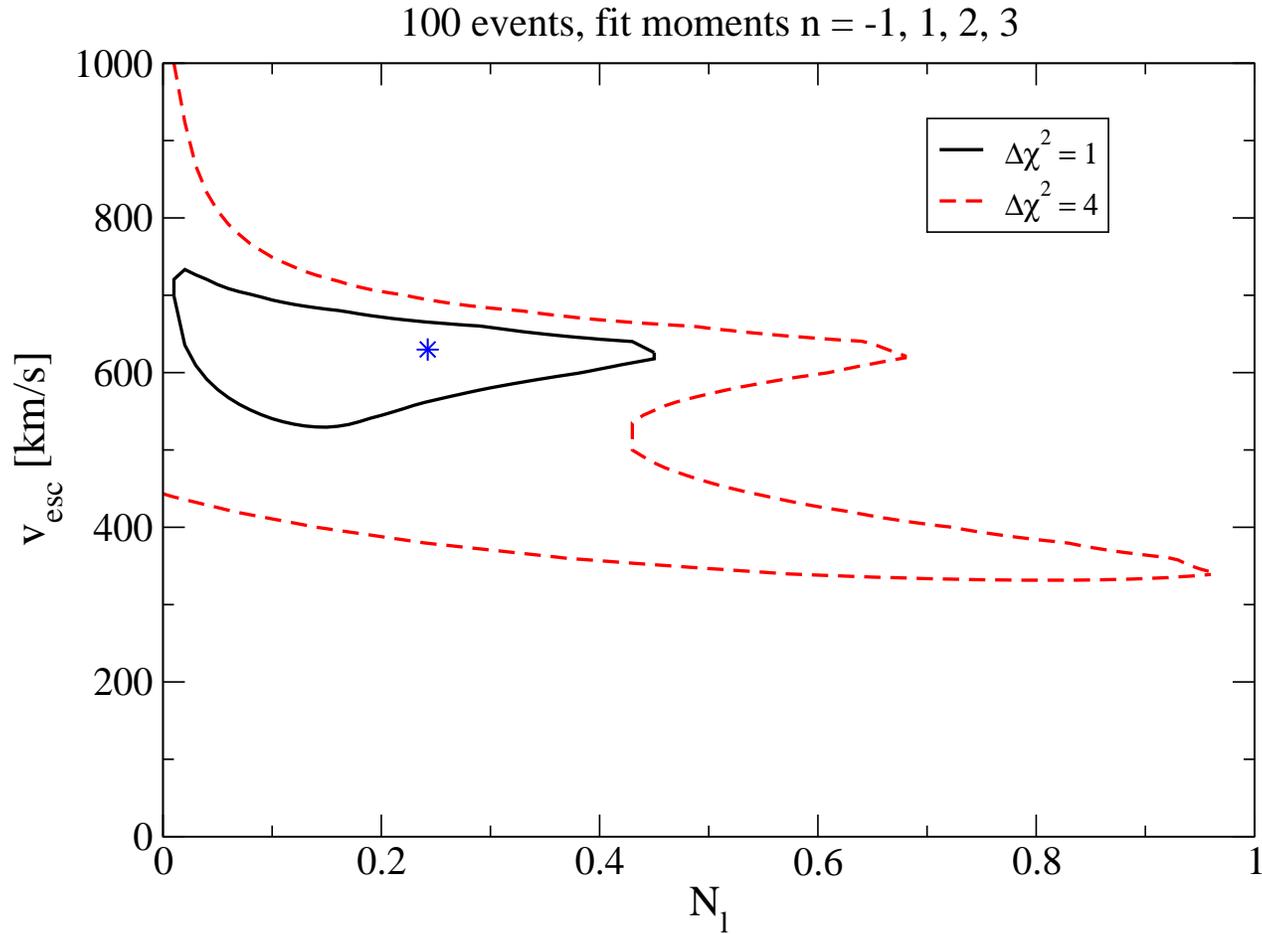
Determination of first 10 moments



Constraining a “late infall” component



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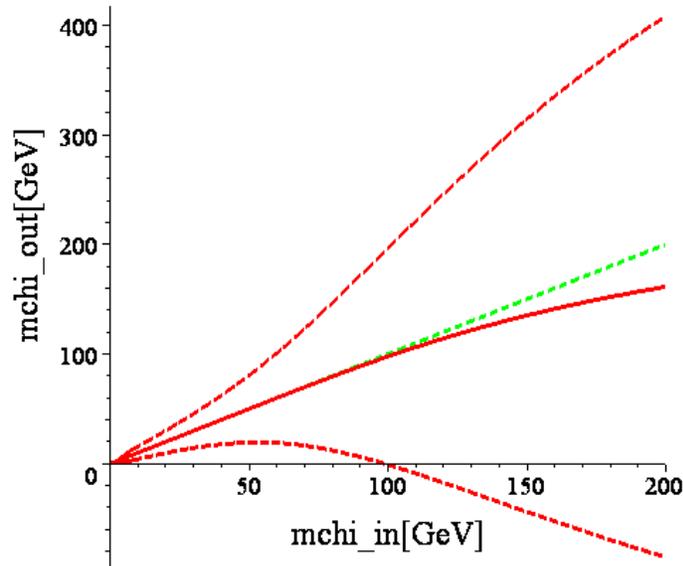


Determining the WIMP mass

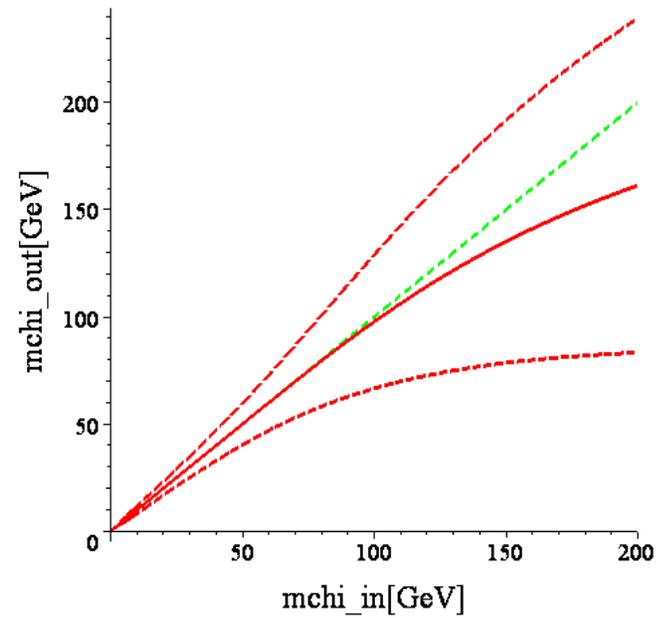
MD & C.L. Shan, in progress

Can determine m_χ from requirement that different targets yield *same* moments of f_1

$Q_{\max} = 200$ keV, $Q_{\min} = 1$ keV, $n = 1, 25 + 25$ events, Ge-76 + Si-28

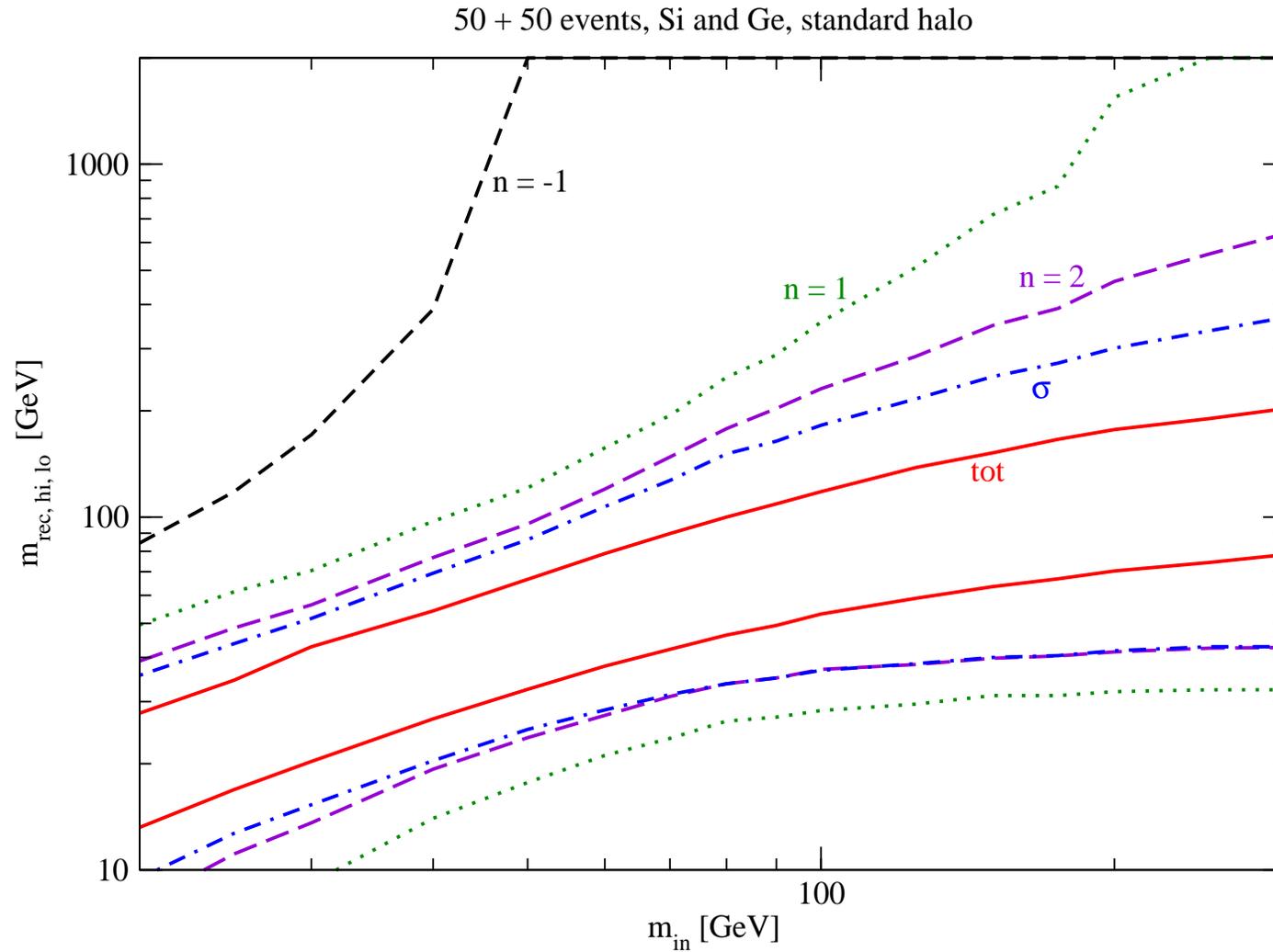


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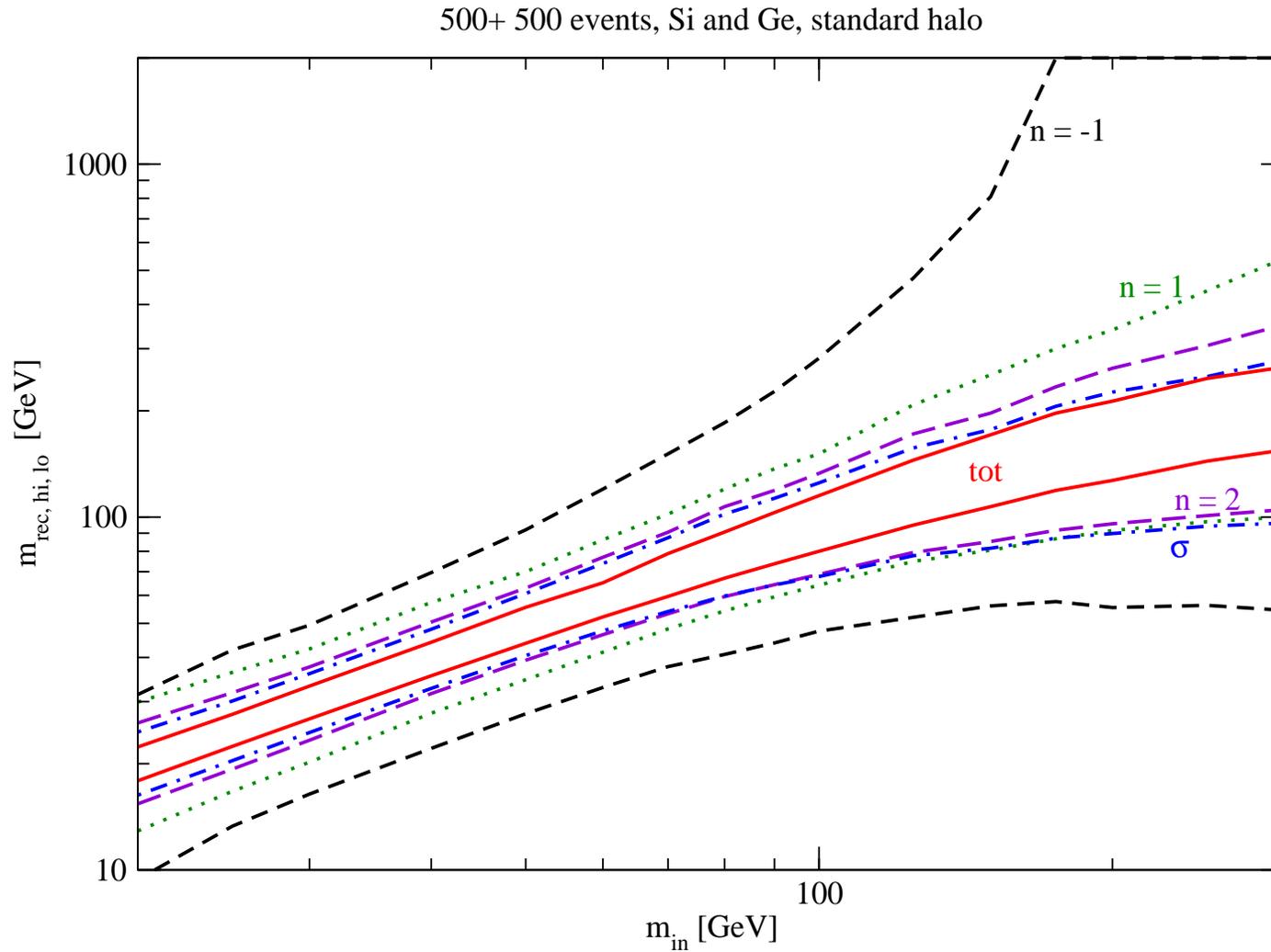
Range of WIMP mass from simulation

Preliminary!



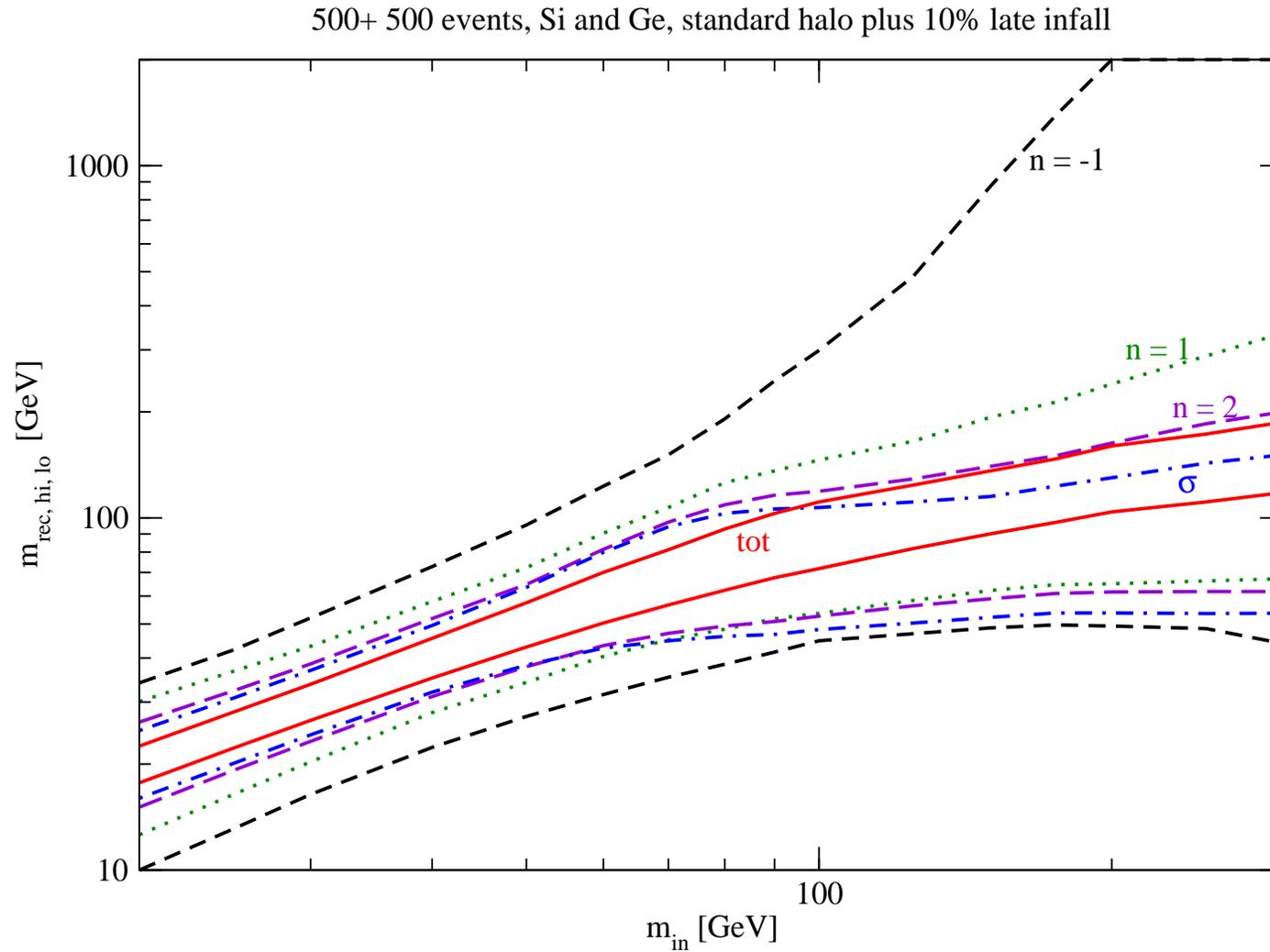
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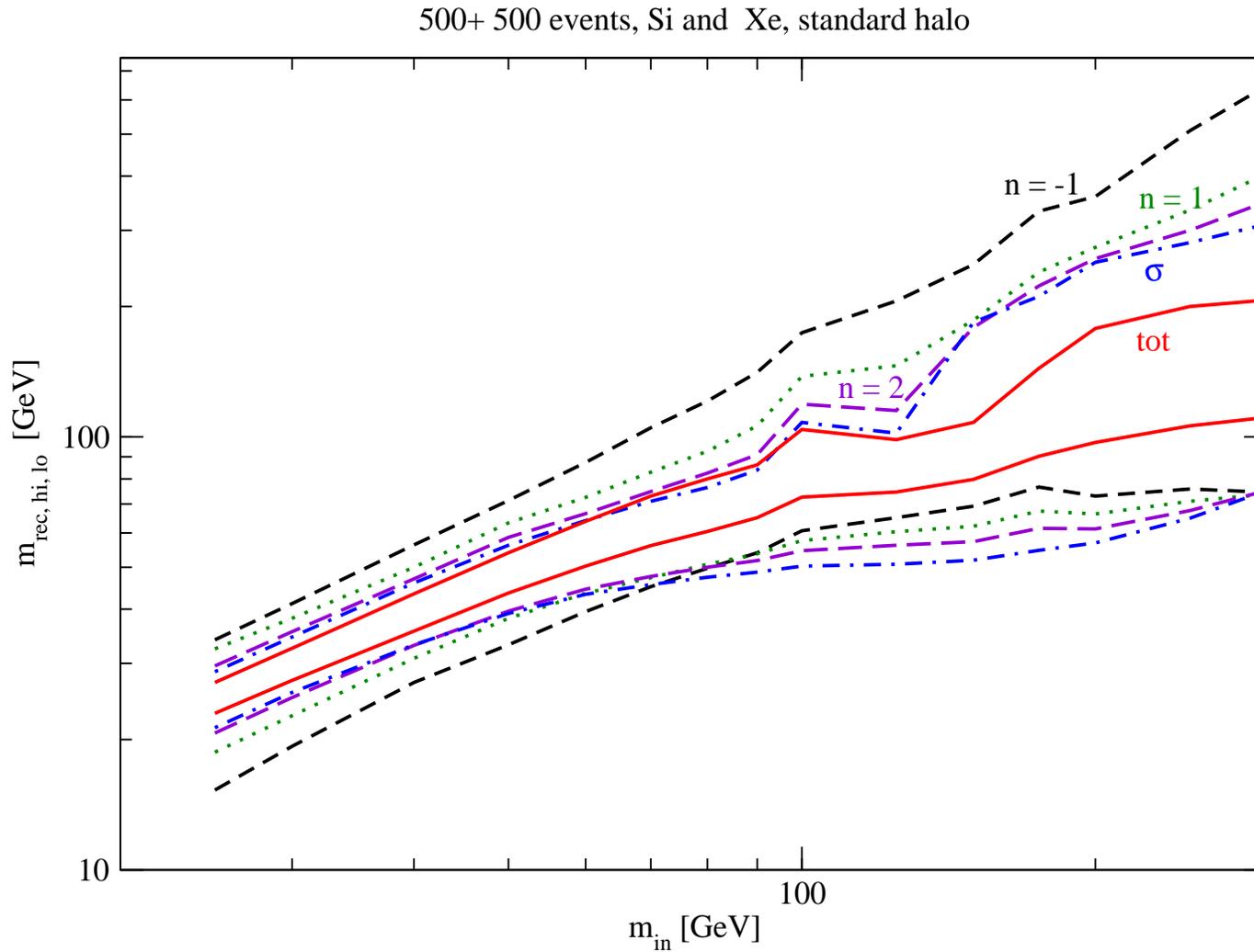
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- **Learning about WIMPs:** Can determine m_χ from moments of f_1 measured with two different targets.