

Exercises on Conformal Field Theory

Dr. Hans Jockers, Andreas Gerhardus

<http://www.th.physik.uni-bonn.de/klemm/cftss16/>

–HOME EXERCISES–

Due on April 29th, 2016

H 2.1 Correlation functions of quasi-primary fields (7 points)

Let us consider correlation function of quasi-primary fields $\phi_k(\vec{x})$ with scaling dimension Δ_k in d dimensions ($d > 2$). We consider conformal invariant ratios in \mathbb{R}^d with $d > 2$.

a) Show that the three-point correlations function is given by

$$\langle \phi_1(\vec{x}_1) \phi_2(\vec{x}_2) \phi_3(\vec{x}_3) \rangle = \frac{C_{123}}{|\vec{x}_1 - \vec{x}_2|^{\Delta_1 + \Delta_2 - \Delta_3} |\vec{x}_1 - \vec{x}_3|^{\Delta_1 + \Delta_3 - \Delta_2} |\vec{x}_2 - \vec{x}_3|^{\Delta_2 + \Delta_3 - \Delta_1}} , \quad (1)$$

with an undetermined real constant C_{123} .

b) Show that the four-point correlation function takes the form

$$\langle \phi_1(\vec{x}_1) \phi_2(\vec{x}_2) \phi_3(\vec{x}_3) \phi_4(\vec{x}_4) \rangle = f(A_1, A_2) \prod_{1 \leq k < l \leq 4} |\vec{x}_k - \vec{x}_l|^{\Delta/3 - \Delta_k - \Delta_l} , \quad (2)$$

where $\Delta = \Delta_1 + \dots + \Delta_4$ and f is an undetermined function of the two anharmonic ratios

$$A_1 = \frac{|\vec{x}_1 - \vec{x}_2| \cdot |\vec{x}_3 - \vec{x}_4|}{|\vec{x}_1 - \vec{x}_3| \cdot |\vec{x}_2 - \vec{x}_4|} , \quad A_2 = \frac{|\vec{x}_1 - \vec{x}_2| \cdot |\vec{x}_3 - \vec{x}_4|}{|\vec{x}_1 - \vec{x}_4| \cdot |\vec{x}_2 - \vec{x}_3|} . \quad (3)$$

Note: The representation of the four-point function given on the right-hand side of eq. (2) is not unique. One could, for example, replace the function f by f/A_1 and absorb the factor A_1 into the product of $|\vec{x}_k - \vec{x}_l|$, thus changing their exponent.

H 2.2 Special conformal transformations (3 points)

The action of the conformal group on a field $\phi(x)$ is defined by specifying the eigenvalues of $\phi(0)$ under $L_{\mu\nu}$, D and K_μ , which generate Lorentz transformations, dilatations and special conformal transformations, respectively:

$$L_{\mu\nu} \phi(0) = S_{\mu\nu} \phi(0) , \quad D \phi(0) = -i\Delta \phi(0) , \quad K_\mu \phi(0) = -\kappa_\mu \phi(0) . \quad (4)$$

Now we use the generator of translations, P_μ , to write $\phi(x) = \exp(-ix \cdot P) \phi(0)$. Together with the conformal algebra

$$\begin{aligned} [D, P_\mu] &= iP_\mu , & [D, K_\mu] &= -iK_\mu , & [K_\mu, P_\nu] &= 2i(\eta_{\mu\nu}D - L_{\mu\nu}) , \\ [K_\rho, L_{\mu\nu}] &= i(\eta_{\rho\mu}K_\nu - \eta_{\rho\nu}K_\mu) , & [P_\rho, L_{\mu\nu}] &= i(\eta_{\rho\mu}P_\nu - \eta_{\rho\nu}P_\mu) , \\ [L_{\mu\nu}, L_{\rho\sigma}] &= i(\eta_{\nu\rho}L_{\mu\sigma} + \eta_{\mu\sigma}L_{\nu\rho} - \eta_{\mu\rho}L_{\nu\sigma} - \eta_{\nu\sigma}L_{\mu\rho}) , \end{aligned} \quad (5)$$

equation (4) thus defines the action of $L_{\mu\nu}$, D and K_μ on $\phi(x)$.

a) Use the Baker–Campbell–Hausdorff formula,

$$e^{-A} B e^A = B + [B, A] + \frac{1}{2!} [[B, A], A] + \frac{1}{3!} [[[B, A], A], A] + \dots, \quad (6)$$

to deduce that

$$K_\mu \phi(x) = [-\kappa_\mu - 2ix_\mu \Delta - 2x^\nu S_{\mu\nu} + 2x_\mu (x \cdot P) - x^2 P_\mu] \phi(x). \quad (7)$$