

Exercises in Superstring Theory

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1 Operator Product Expansions

In a 2d conformal invariant theory, there are two non-vanishing components of the energy-momentum tensor: the chiral part $T(z)$ and the anti-chiral part $\bar{T}(\bar{z})$. They give rise to an infinite number of conserved currents from which the associated conserved charges generating infinitesimal conformal transformations $z \rightarrow z + \epsilon(z)$, $\epsilon(z) \ll 1$, are written as

$$Q_\epsilon = \oint_{C_0} \frac{dz}{2\pi i} \epsilon(z) T(z) . \quad (1.1)$$

Any field $\phi(z)$ on the complex plane has a mode expansion of the following form

$$\phi(z) = \sum_{n \in \mathbb{Z}} \phi_n z^{-n-h} \quad \text{with} \quad \phi_n = \oint_{C_0} \frac{dz}{2\pi i} \phi(z) z^{n+h-1}, \quad (1.2)$$

where integration is counterclockwise around the origin,

1. For two generic operators $A(z)$ and $B(z)$ with

$$A = \oint_{C_0} \frac{dz}{2\pi i} A(z) , \quad B = \oint_{C_0} \frac{dz}{2\pi i} B(z) , \quad (1.3)$$

show that radial ordering \mathcal{R} implies that

$$\oint_{C_0} \frac{dz}{2\pi i} [A(z), B(w)] = \oint_{C_w} \frac{dz}{2\pi i} \mathcal{R}(A(z)B(w)) . \quad (1.4)$$

In particular, this leads to the important relation

$$[A, B] = \oint_{C_0} \frac{dw}{2\pi i} \oint_{C_w} \frac{dz}{2\pi i} \mathcal{R}(A(z)B(w)) . \quad (1.5)$$

2. The variation (2.2) of exercise sheet 6 is encoded into the commutator $\delta_\epsilon \phi(w) = -[Q_\epsilon, \phi(w)]$. Use the result from item 1. and the Cauchy-Riemann formula given by

$$\oint_{C_w} \frac{dz}{2\pi i} \frac{f(z)}{(z-w)^n} = \frac{1}{(n-1)!} f^{(n-1)}(w) \quad (1.6)$$

to show that the (\mathcal{R} -ordered) *operator product expansion* $T(z)\phi(w)$ is given by

$$T(z)\phi(w) = \frac{h\phi(w)}{(z-w)^2} + \frac{\partial_w \phi(w)}{z-w} + \text{reg.} , \quad (1.7)$$

where reg. stands for a holomorphic fct. of z regular at $z = w$, also called finite terms or non-singular terms.

Therefore, a primary field $\phi(z)$ of weight h is also defined as a field which has the operator product expansion (1.7) with the energy-momentum tensor.

Consider the Laurent expansion of the energy-momentum tensor $T(z)$ given by

$$T(z) = \sum_{n \in \mathbb{Z}} z^{-n-2} L_n, \quad (1.8)$$

where the L_n 's are the Virasoro generators given by

$$L_n = \oint \frac{dz}{2\pi i} z^{n+1} T(z). \quad (1.9)$$

3. Show that the (radial-ordered) OPE of the energy-momentum tensor with itself,

$$T(z)T(w) = \frac{c/2}{(z-w)^4} + \frac{2T(w)}{(z-w)^2} + \frac{\partial_w T(w)}{(z-w)} + \text{finite terms}, \quad (1.10)$$

is equivalent to the Virasoro algebra

$$[L_m, L_n] = (m-n)L_{m+n} + \frac{c}{12}m(m^2-1)\delta_{m+n,0}. \quad (1.11)$$

2 The propagator of the free boson

The first simple system is that of a free boson. In two dimensions the free boson has the following Euclidean action

$$S = \frac{g}{2} \int d^2x (\partial_\alpha \varphi \partial^\alpha \varphi + m^2 \varphi^2), \quad (2.1)$$

where g is a normalization parameter. The two-point function called *propagator* is given by

$$K(\mathbf{x}, \mathbf{y}) = \langle \varphi(\mathbf{x}) \varphi(\mathbf{y}) \rangle. \quad (2.2)$$

The propagator obeys

$$g(-\partial_x^2 + m^2)K(\mathbf{x}, \mathbf{y}) = \delta(\mathbf{x} - \mathbf{y}). \quad (2.3)$$

1. Why is the propagator a function of only $r := |\mathbf{x} - \mathbf{y}|$, i.e. $K(\mathbf{x}, \mathbf{y}) = K(r)$? Show that for $m = 0$ (2.3) is solved by

$$K(r) = -\frac{1}{2\pi g} \log r + \text{const}. \quad (2.4)$$

Hint: Proceed by integrating (2.3) over \mathbf{x} within a disc D of radius r centered around \mathbf{y}

In terms of complex coordinates this reads

$$\langle \varphi(z, \bar{z}) \varphi(w, \bar{w}) \rangle = -\frac{1}{4\pi g} \left(\log(z-w) + \log(\bar{z}-\bar{w}) \right) + \text{const}. \quad (2.5)$$

This leads to the following OPE

$$\partial_z \varphi(z, \bar{z}) \partial_w \varphi(w, \bar{w}) \sim -\frac{1}{4\pi g} \frac{1}{(z-w)^2}. \quad (2.6)$$

2. Given the energy momentum tensor

$$T(z) = -2\pi g : \partial \varphi \partial \varphi : \quad (2.7)$$

Show that the normal ordered operators $V_\alpha(z, \bar{z}) =: e^{i\alpha\varphi(z, \bar{z})}$ are primary fields and determine their conformal weights h and \bar{h} .

Hint: Determine the OPE with the energy momentum tensor $T(z)$.

3 The Free Fermion

The action for a free Majorana fermion reads

$$S = \frac{1}{2}g \int d^2x \Psi^\dagger \gamma^0 \gamma^\mu \partial_\mu \Psi, \quad (3.1)$$

where the Dirac matrices γ^μ satisfy the so-called Dirac algebra

$$\gamma^\mu \gamma^\nu + \gamma^\nu \gamma^\mu = 2\eta^{\mu\nu}, \quad (3.2)$$

if $\eta^{\mu\nu} = \text{diag}(1, 1)$ a representation thereof is

$$\gamma^0 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \gamma^1 = i \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \quad (3.3)$$

1. Write the action in S in terms of the two-component spinor $\Psi = (\psi, \bar{\psi})$ and calculate the equations of motion for ψ and $\bar{\psi}$. What do they imply?
2. Next we want to calculate the correlator $\langle \Psi_i(\mathbf{x}), \Psi_j(\mathbf{y}) \rangle$, where $i, j = 1, 2$ label the components of Ψ . To do so, express the kinetic terms in the derived action in 1. in terms of a matrix A_{ij} and write a differential equation for the Green's function.
3. We claim the the Green's function $G_{ij}(z, \bar{z})$ in complex coordinates for item 2. is given by

$$G = \frac{1}{2\pi g} \begin{pmatrix} \frac{1}{z-w} & 0 \\ 0 & \frac{1}{\bar{z}-\bar{w}} \end{pmatrix}. \quad (3.4)$$

Prove this by using the techniques you used in the bosonic case.

4 The Ghost System

Another simple system is the so-called *ghost system* with the following action

$$S = \frac{g}{2} \int d^2x b_{\mu\nu} \partial^\mu c^\nu, \quad (4.1)$$

where the field $b_{\mu\nu}$ is a traceless symmetric tensor, and both $b_{\mu\nu}$ and c^μ are anticommuting fields. The propagator of such system can be obtained in a similar way as sketched in section 2. It is given by

$$\langle b(z)c(w) \rangle = \frac{1}{\pi g} \frac{1}{z-w}. \quad (4.2)$$

1. Determine the correlators $\langle b(z)\partial c(w) \rangle$, $\langle \partial b(z)c(w) \rangle$ and $\langle \partial b(z)\partial c(w) \rangle$.
2. The normal-ordered holomorphic energy-momentum tensor of the bc -system is given by

$$T(z) = \pi g : (2\partial cb + c\partial b) : \quad (4.3)$$

Compute the OPEs $T(z)b(w)$ and $T(z)c(w)$ using Wick's Theorem. What are the conformal dimensions of the fields b and c ?

3. Compute the OPE of the energy-momentum tensor with itself and bring it to the form

$$T(z)T(w) = \frac{c/2}{(z-w)^4} + \frac{hT(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w} + \text{ref}. \quad (4.4)$$

Read off the conformal weight h of the energy momentum tensor and the central charge c of the bc system.