Exercise 9 14th June 2013 SS 13

8 points

## **Exercises on Theoretical Particle Astrophysics**

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## 9.1 Dark Matter in Galaxies

The observation of the galactic rotation curve<sup>1</sup> 1 yields a deficit of mass in the galaxy. Under the assumption of spherical symmetry of a rotating galaxy one can calculate the mass inside a sphere of a given radius from the circular velocity of the stars at its surface and compare it to an estimation from the visible stars.



## Observed vs. Predicted Keplerian

Figure 1: rotation curve of a galaxy

- (a) Give a formula which expresses the circular velocity in terms of the enclosed mass and the distance to the galactic center. Verify the virial theorem for gravitationally bound systems  $\langle T \rangle = -\langle V \rangle /2$ . (2 points)
- (b) Assume the simplest case of a constant mass density  $\rho_0$  inside a radius  $r_0$ . How does the rotation curve look like? (1 point)

<sup>&</sup>lt;sup>1</sup>Image taken from *http*://www.astronomy.ohio-state.edu/thompson/162/Lecture40.html

(c) A more realistic distribution is of the form

$$\rho(r) = \frac{\rho_0 r_0^2}{r^2 \left(1 + r/r_0\right)^{\alpha}}$$

Derive the rotation curve v(r). Which value of  $\alpha$  gives a flat rotation curve at  $r \gg r_0$  as shown in the measurements? (3 points)

(d) At  $r = 10^5$  light years the measurement yields  $v_{calc} = 15$  km/s and  $v_{meas} = 225$  km/s. Calculate the visible as well as the true galaxy mass. What is the percentage of dark matter in the galaxy? How high is the average dark matter mass density? Hint:  $G = 6.67 \times 10^{-11} m^3 kg^{-1} s^{-2}$  (2 points)

## 9.2 The Boltzmann equation

Consider a stable particle  $\psi$ . In a comoving volume, we know that the number of  $\psi$  and  $\overline{\psi}$  changes only through annihilation and inverse annihilation processes (with  $\chi$  we indicate all the possible final states):

$$\psi \overline{\psi} \leftrightarrow \chi \overline{\chi}$$
.

Under certain simplifying assumptions, the Boltzmann equation that rules the evolution of the number density for  $\psi$  and can be written as:

$$\frac{dn_{\psi}}{dt} + 3Hn_{\psi} = -\langle \sigma_A | v | \rangle (n_{\psi}^2 - (n_{\psi}^{\mathrm{EQ}})^2) , \qquad (1)$$

where  $\sigma_A|v|$  is the total annihilation cross section, and  $n_{\psi}^{\text{EQ}}$  is the particle number density at thermal equilibrium. Let us take a system in which the assumptions that lead to the previous formula are fulfilled, and consider the following questions:

(a) Take a particle  $\psi$  and use the following quantity

$$Y = \frac{n_{\psi}}{s}$$

where s is the entropy density. Using the conservation of entropy per comoving volume  $(sa^3 = \text{constant})$ , show that (1 point)

$$\dot{n}_{\psi} + 3Hn_{\psi} = s\dot{Y}.$$
(2)

(b) Let m be the mass of the particle  $\psi$ . Now introduce the quantity

$$x \equiv \frac{m}{T} \,. \tag{3}$$

During the radiation dominated era, define also  $H(m) \simeq 1.67(g^*)^{\frac{1}{2}}m^2/m_{Pl}$ , and  $H(x) = H(m)x^{-2}$ . Show that the Boltzmann equation becomes (4 points)

$$\frac{dY}{dx} = \frac{-x - \langle \sigma_A | v | \rangle s}{H(m)} (Y^2 - Y_{\rm E}^2) \,. \tag{4}$$

12 points

- (c) Write the expression for  $Y_{EQ}(x)$  (notice, as a function of x), in the case  $x \gg 3$  (that is, the non-relativistic limit), and in the case  $3 \gg x$  (the relativistic limit). Suppose that the freezed out occurs at  $x \equiv x_f$  while still in the relativistic case. What is the value of  $Y_{EQ}(x)$  at  $x_f$ ? (3 points)
- (d) We have derived the x-dependent Boltzmann equation

$$\frac{dY}{dx} = \frac{\lambda}{x^2} (Y^2 - Y_{\rm E}^2), \qquad (5)$$

where  $\lambda$  is parametrized by

$$\lambda = \frac{m^3 \langle \sigma_A | v | \rangle}{H(m)} \tag{6}$$

and can be considered as constant in this exercise. At late times, i.e. well after freezeout, Y will be much larger than  $Y_{\rm EQ}$  and the relation

$$\frac{dY}{dx} \simeq \frac{\lambda Y^2}{x^2} (x \ll 1) \tag{7}$$

holds. Integrate equation (7) analytically in order to derive the approximation

$$Y_{\infty} \simeq \frac{x_f}{\lambda}$$
 (8)

Typically one can consider  $Y_f$  being significantly larger than  $Y_{\infty}$ . (4 points)