

General Relativity and Cosmology

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Example sheet 10

1. Motion in Schwarzschild solution

Consider the Schwarzschild metric written in the general form (see ex. sheet 5, problem 4)

$$ds^2 = A(r) dt^2 + B(r) dr^2 + r^2 [d\theta^2 + \sin^2 \theta d\phi^2] \quad (1)$$

(a) Write the equations of motion for this metric (keep for the moment A and B general)

(b) Setting $\theta = \pi/2$ (why can this be done?), integrate suitably to get

$$\frac{dt}{dp} = \frac{1}{B(r)}, \quad r^2 \frac{d\phi}{dp} = J = \text{const.} \quad (2)$$

and

$$B(r) \left(\frac{dr}{dp} \right)^2 + \frac{J^2}{r^2} - \frac{1}{A(r)} = -E = \text{const.} \quad (3)$$

where p is the parameter along the worldline.

(c) Show that $d\tau^2 = E dp^2$, and therefore $E = 0$ must hold for photons, while $E > 0$ for other matter.

(d) Eliminate dp from the integrals of motion obtained in b) to get a relation between r and ϕ . Show that

$$\phi = \pm \int \frac{\sqrt{B(r)} dr}{r^2 \sqrt{\frac{1}{J^2 A(r)} - \frac{E}{J^2} - \frac{1}{r^2}}} \quad (4)$$

is a solution.

2. Light deflection

A photon approaches the central mass from infinity along the direction $\phi_\infty = 0$ with impact parameter b . We want to calculate the deflection of its trajectory. Let r_0 be the radius of the closest approach.

(a) Determine the value of J in terms of r_0

(b) Show that now (4) reduces to

$$\phi(r) = \phi_\infty + \int_r^\infty \frac{\sqrt{B(r')}}{r'^2 \sqrt{\frac{r'^2 A(r_0)}{r_0^2 A(r')} - 1}} \frac{dr'}{r'} \quad (5)$$

(c) Show that the Schwarzschild line element calculated in the class, can be approximated by

$$B(r) \sim 1 + \frac{2MG}{r}, \quad A(r) \sim 1 - \frac{2MG}{r} \quad (6)$$

in regions where Newtonian gravity is valid.

(d) Use (5) to calculate the deflection angle $\Delta\phi = 2|\phi(r_0) - \phi_\infty| - \pi$.

Use this equality which holds to lowest order in $2MG/r$:

$$\frac{r^2}{r_0^2} \frac{A(r_0)}{A(r)} - 1 = \left[\frac{r^2}{r_0^2} - 1 \right] \left[1 - \frac{2MG}{r_0(r+r_0)} \right] \quad (7)$$

The following integrals may be useful:

$$\int \frac{dx}{x\sqrt{x^2 - a^2}} = \frac{1}{a} \arccos \frac{a}{x}, \quad \int \frac{dx}{x^2\sqrt{x^2 - a^2}} = \frac{\sqrt{x^2 - a^2}}{a^2 x}, \quad \int \frac{dx}{(x+a)\sqrt{x^2 - a^2}} = \frac{\sqrt{x^2 - a^2}}{a(x+a)}.$$

3. Precession of perihelia

Use eqs. (3, 4, 6) to

(a) Determine E and J^2 by looking at the aphelion $r = r_+$ and perihelion $r = r_-$ of a planet in bound orbit around the sun, for general $A(r)$. (*At r_\pm , $dr/d\phi$ vanishes.*)

(b) Show that the amount of orbital precession per revolution is

$$\Delta\phi = 2|\phi(r_+) - \phi(r_-)| - 2\pi, \quad (8)$$

where

$$\phi(r_+) - \phi(r_-) = \int_{r_-}^{r_+} \left[\frac{r_-^2 (A^{-1}(r) - A^{-1}(r_-)) - r_+^2 (A^{-1}(r) - A^{-1}(r_+))}{r_-^2 r_+^2 (A^{-1}(r_+) - A^{-1}(r_-))} - \frac{1}{r^2} \right]^{-1/2} \times \frac{\sqrt{B(r)} dr}{r^2} \quad (9)$$

(c) Show that for weak fields, we can use

$$A^{-1}(r) = 1 + \frac{2MG}{r} + \frac{4M^2G^2}{r^2} \quad (10)$$

which makes the first term in (9) quadratic in $1/r$, and that we can then write

$$\phi(r_+) - \phi(r_-) = \int_{r_-}^{r_+} \left[C \left(\frac{1}{r_-} - \frac{1}{r} \right) \left(\frac{1}{r} - \frac{1}{r_+} \right) \right]^{-1/2} \times \frac{\sqrt{B(r)} dr}{r^2} \quad (11)$$

(d) Determine C in the limit $r \rightarrow \infty$. you should get

$$C \sim 1 - \frac{4MG}{L} + \dots, \quad (12)$$

where

$$\frac{1}{L} = \frac{1}{2} \left(\frac{1}{r_+} + \frac{1}{r_-} \right).$$

Show that (9) now reduces to

$$\phi(r_+) - \phi(r_-) = \left(1 + \frac{2MG}{L} \right) \times \int_{r_-}^{r_+} \frac{\left(1 + \frac{MG}{r} \right) dr}{r^2 \sqrt{\left(\frac{1}{r_-} - \frac{1}{r} \right) \left(\frac{1}{r} - \frac{1}{r_+} \right)}} \quad (13)$$

(e) Calculate $\Delta\phi$. (*You can approximate the result of the integral with $(1 + \frac{MG}{L}\pi)$.*)

(f) Determine the total precession $\Delta\phi$ for Mercury over the time of a century. (415 revolutions per century; $L = 55.3 \times 10^9\text{m}$; $MG = 1475\text{m}$). The observed value is 43.11 ± 0.45 arcseconds.