

## Exercises on Theoretical Particle Physics II

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### 4.1 Vectorsuperfields and gauge transformations

(10 credits)

Global  $\mathcal{N} = 1$  supersymmetry allows one more supermultiplet, the **vectormultiplet**, which contains a supersymmetric version of a gauge theory. It consists of the usual spin one gauge boson  $V_\mu$  as well as its spin one half superpartner  $\lambda$  called the **gaugino**. There also exists a superfield formulation of the vectormultiplet completely analogous to the chiral superfield describing the chiral multiplet  $(\varphi, \psi)$ . The appropriate superfield  $V$  is the **vectorsuperfield** defined by  $V = V^\dagger$  with the expansion

$$\begin{aligned}
 V(x, \theta, \bar{\theta}) = & C(x) + i\theta\chi(x) - i\bar{\theta}\bar{\chi}(x) - \theta\sigma^\mu\bar{\theta}V_\mu(x) \\
 & + \frac{1}{2}i\theta\theta[M(x) + iN(x)] - \frac{1}{2}i\bar{\theta}\bar{\theta}[M(x) - iN(x)] \\
 & + i\theta\theta\bar{\theta}\left[\bar{\lambda}(x) + \frac{i}{2}\bar{\sigma}^\mu\partial_\mu\chi(x)\right] - i\bar{\theta}\bar{\theta}\theta\left[\lambda(x) + \frac{i}{2}\sigma^\mu\partial_\mu\bar{\chi}(x)\right] \\
 & + \frac{1}{2}\theta\theta\bar{\theta}\bar{\theta}\left[D(x) + \frac{1}{2}\partial_\mu\partial^\mu C(x)\right].
 \end{aligned} \tag{1}$$

- (a) Check that (1) is indeed a vectorsuperfield. (3 credits)
- (b) Compare the expansion of  $V$  to the one of the vectorsuperfield defined by  $\Lambda + \Lambda^\dagger$  where  $\Lambda_L(x, \theta) = \Lambda(x) + \sqrt{2}\theta\psi_\Lambda(x) + \theta\theta F_\Lambda(x)$  is a left-chiral superfield, here given in the left-chiral representation. What is the interpretation of the transformation

$$V \mapsto V + \Lambda + \Lambda^\dagger ? \tag{2}$$

Write the transformation law for all the component fields of  $V$ . (4 credits)

*Hint: Work in the left-chiral representation by shifting the argument  $x^\mu$  of  $\Lambda^\dagger$ . How does  $V_\mu$  transform?*

- (c) Give the appropriate  $\Lambda$  in (2) to transform  $V$  into the **Wess-Zumino gauge**  $V_{\text{WZ}}$ , i.e. to obtain  $C(x) = \chi(x) = M(x) = N(x) = 0$ . What is the highest non-vanishing power of  $V_{\text{WZ}}$ ? Calculate  $V_{\text{WZ}}$ ,  $V_{\text{WZ}}^2$  as well as  $V_{\text{WZ}}^3$ . (3 credits)

## 4.2 Supersymmetric U(1) gauge theory.

(10 credits)

To construct an action exhibiting the gauge symmetry (2) we have to find an adequate gauge invariant quantity. This will be the building block of any gauge invariant action. Exploiting the gauge invariance of the gaugino  $\lambda$  w.r.t. (2) in WZ-gauge we define the supersymmetric field strength of  $V$  by

$$W_\alpha = -\frac{1}{4}\bar{D}\bar{D}D_\alpha V, \quad \bar{W}_{\dot{\alpha}} = -\frac{1}{4}DD\bar{D}_{\dot{\alpha}}V. \quad (3)$$

Note that the lowest component (in  $\theta, \bar{\theta}$ ) of  $W_\alpha$  is the gauge invariant gaugino  $\lambda_\alpha$ .

- (a) Show that  $W_\alpha$  in (3) defines a gauge invariant, left-chiral superfield! How do  $W_\alpha, \bar{W}_{\dot{\alpha}}$  transform under Lorentz transformations? Expand  $W_\alpha$  in its component fields. (5 credits)

*Hint: Translate  $V_{WZ}, D_\alpha$  and  $\bar{D}_{\dot{\alpha}}$  into the left-chiral representation. Finally use  $\sigma^\mu\bar{\sigma}^\nu - \eta^{\mu\nu} = -2i\sigma^{\mu\nu}$ .*

- (b) The simplest SUSY, gauge as well as Lorentz-invariant action for a vectorsuperfield reads

$$S_{U(1)} = \int d^4x d\theta^2 W^\alpha W_\alpha. \quad (4)$$

Why is this SUSY-invariant? Determine its component expression

$$\mathcal{L}_{U(1)} = -\frac{1}{2}F^{\mu\nu}F_{\mu\nu} - 2i\lambda\sigma^\mu\partial_\mu\bar{\lambda} + D^2 + \frac{i}{4}\epsilon^{\mu\nu\rho\sigma}F_{\mu\nu}F_{\rho\sigma}, \quad (5)$$

with field strength  $F_{\mu\nu}$ . The last term is imaginary and cancels after adding the h.c. to (4). Note also the presence of the new auxiliary field  $D$ . (5 credits)

*Hint: Use the identities  $\text{tr}(\sigma^{\mu\nu}) = 0$  and  $\sigma^{\mu\nu}\sigma^{\rho\sigma} = \frac{1}{4}(\eta^{\mu\rho}\eta^{\nu\sigma} - \eta^{\mu\sigma}\eta^{\nu\rho}) - \frac{i}{4}\epsilon^{\mu\nu\rho\sigma}$ .*