
Exercises on General Relativity and Cosmology

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–EXERCISE ONE–

Due on April 10

H1 Field Theory: The Example of Classical Electrodynamics (15 points)

In this exercise we consider the field theoretical formulation of classical electrodynamics. The electromagnetic field is described in terms of a vector field

$$A : \mathbb{R}^{1,3} \longrightarrow \mathbb{R}^{1,3}, \quad (1)$$

with components A^μ . From these components, we can build the field strength tensor F , via

$$F_{\mu\nu} = \frac{\partial}{\partial x^\mu} A_\nu - \frac{\partial}{\partial x^\nu} A_\mu. \quad (2)$$

The particles of mass m_i and charge q_i are described by their trajectories,

$$x_i : I_k \subset \mathbb{R} \longrightarrow \mathbb{R}^{1,3} \quad \text{for } i = 1, \dots, N = \text{number of particles} \quad (3)$$

$$\sigma_i \longmapsto x_i(\sigma_i), \quad (4)$$

For each particle, then σ is an arbitrary parameter used to parameterize its trajectory, a specific point on the trajectory curve then corresponds to coordinates x^μ . Written in coordinates, the total action is

$$S[x_i, A] = - \sum_{i=1}^N \int_{I_k} d\sigma_i \left(m_i \sqrt{-\eta_{\alpha\beta} \dot{x}_i^\alpha(\sigma_i) \dot{x}_i^\beta(\sigma_i)} - q_i A_\alpha(x_i(\sigma_i)) \dot{x}_i^\alpha \right) \quad (5)$$

$$- \frac{1}{4} \int_{\mathbb{R}^{1,3}} d^4x F_{\alpha\beta}(x) F^{\alpha\beta}(x), \quad (6)$$

where $\dot{x}_i = \frac{d}{d\sigma_i} x_i$ and integration on Minkowski space is the same as on \mathbb{R}^4 .

a) Shortly comment on the significance of each term in S . (1 point)

b) Take the variation of S with respect to x_i^μ in order to derive the Einstein-Lorentz equation (2 points)

$$m_i \frac{d}{d\sigma_i} \frac{\dot{x}_i^\mu(\sigma_i)}{\sqrt{-\eta_{\alpha\beta} \dot{x}_i^\alpha(\sigma_i) \dot{x}_i^\beta(\sigma_i)}} = q_i F^\mu{}_\nu(x_i(\sigma_i)) \dot{x}_i^\nu. \quad (7)$$

- c) Rewrite the second term in S , the term where q_i appears, in terms of the charge-current density (1 point)

$$j^\mu(x) = \sum_{i=1}^N q_i \int d\sigma_i \delta^{(4)}(x - x_i(\sigma_i)) \dot{x}_i^\mu(\sigma_i). \quad (8)$$

- d) Take the variation of S with respect to A_μ to derive the inhomogenous Maxwell's equations (2 points)

$$\frac{\partial}{\partial x^\mu} F^{\nu\mu}(x) = j^\nu(x). \quad (9)$$

- e) Use the definition of the field strength tensor, eq. (2), to show the homogenous Maxwell's equations (2 points)

$$\frac{\partial}{\partial x^\alpha} F_{\mu\nu} + \frac{\partial}{\partial x^\mu} F_{\nu\alpha} + \frac{\partial}{\partial x^\nu} F_{\alpha\mu} = 0. \quad (10)$$

Now take $A^\mu = (\phi, \vec{A})$ with ϕ and \vec{A} such that $\vec{B} = \text{rot } \vec{A}$ and $\vec{E} = -\text{grad } \phi - \dot{\vec{A}}$. Further, $j^\mu = (\rho, \vec{j})$ with the charge-density ρ and current-density \vec{j} .

- f) Express the components of the field strength tensor, $F_{\alpha\beta}$, in terms of the components of the electric and magnetic field, \vec{E} and \vec{B} . (1 point)

- g) Show that eq. (9) indeed gives the inhomogenous Maxwell's equations: (2 points)

$$\text{div } \vec{E} = \rho, \quad \text{rot } \vec{B} - \dot{\vec{E}} = \vec{j}. \quad (11)$$

- h) Show that eq. (10) indeed gives the homogenous Maxwell's equations: (2 points)

$$\text{div } \vec{B} = 0, \quad \text{rot } \vec{E} + \dot{\vec{B}} = 0. \quad (12)$$

- i) Parameterise eq. (7) by time, i.e. $\sigma = x^0 = t$, and show that it reduces to (2 points)

$$m \frac{d}{dt} \frac{\vec{v}}{\sqrt{1-v^2}} = q \left(\vec{E} + \vec{v} \times \vec{B} \right). \quad (13)$$