
Exercises on General Relativity and Cosmology

Dr. Stefan Förste, Bardia Najjari

<http://www.th.physik.uni-bonn.de/people/bardia/GRCss19/GR.html>

–EXERCISE FIVE :)–
Due May 8th

H 5.1 Vectors and Vector fields

(12 points)

In the lecture, we introduced the directional derivatives along curves passing through a point P on the manifold, as the tangent space T_P . We would like to first confirm that the directional derivative does indeed form a vector space.

To that end; we should check that if $\frac{d}{d\lambda}$ and $\frac{d}{d\gamma}$ are two directional derivatives along two curves parameterized by λ and γ , then the combined operator $\mathcal{O} = a\frac{d}{d\lambda} + b\frac{d}{d\gamma}$, where a and b are constant numbers, is also gonna be in the same space. That is to say that \mathcal{O} is a directional derivative. For an operator to be a well-defined derivative (*derivation*), it should satisfy

$$\begin{aligned} \text{Linearity:} \quad & X(\alpha f + \beta g) = \alpha X(f) + \beta X(g) \quad \text{with } \alpha, \beta \in \mathbb{R}, f, g \in C^\infty \\ \text{Leibniz rule:} \quad & X(f \cdot g) = f \cdot X(g) + g \cdot X(f) \quad \text{with } f, g \in C^\infty. \end{aligned} \quad (1)$$

a) show by explicit calculation that \mathcal{O} is a derivative. (2 points)

Let us move further with the vectors on a manifold. So far, we were working with the vector space defined at a point P on a manifold; pinning a vector space at each and every point on a manifold, we will have a *vector field*. Just as the vector space at point P was a map from functions to their derivative functions at the point P , a vector field maps functions to derivative functions, all over the manifold. A smooth vector field X on a manifold M fulfils the two above conditions, all over the manifold.

Given two vector fields X and Y we define a new vector field $[X, Y]$, the *Lie bracket* or *commutator* of X and Y , by

$$[X, Y](f) = X(Y(f)) - Y(X(f)) \quad \text{for } f \in C^\infty(M). \quad (2)$$

b) Show in two ways that $[X, Y]$ is indeed a vector field:

- i) Prove that $[X, Y]$ is a derivation.
- ii) Write $[X, Y]$ in terms of *components* and show that they transform as those of a vector field under change of coordinates.

(2 points)

- c) What about XY or YX ? are these combinations of vector fields, also vector fields themselves? (2 points)
- d) Show that the Lie bracket is (2 points)
- i) skew-symmetric, $[X, Y] = -[Y, X]$, and
- ii) satisfies the Jacobi identity, $[[X, Y], Z] + [[Z, X], Y] + [[Y, Z], X] = 0$.
- e) Consider as the manifold, \mathbb{R}^2 equipped with some coordinates x^1, x^2 . Calculate the Lie bracket of the coordinate vector fields $\partial_1 = \frac{\partial}{\partial x^1}$ and $\partial_2 = \frac{\partial}{\partial x^2}$. (1 point)
- f) Find an example of two nowhere-vanishing, (at each point) linearly independent vector fields in \mathbb{R}^2 whose Lie bracket does not vanish. Note that these two vector fields provide a basis for the tangent space at each point. Can you find coordinates, under which this basis will be a coordinate basis? (3 points)

H 5.2 Tensors and (A)symmetry

(4 points)

On the previous sheets, you already used the symmetry/anti-symmetry properties of tensors occasionally. We would like to do get a bit more used to tensors in that respect. Let T be a tensor in the vector space $T^{0,k}(V)$. Given this tensor, one can then define the two related tensors.

$$T_{\mu_1, \dots, \mu_k}^{sym} := T_{(\mu_1, \dots, \mu_k)} = \frac{1}{k!} \sum_{\sigma \in S_K} T_{\mu_{\sigma_1}, \dots, \mu_{\sigma_k}} \quad (3)$$

is the corresponding symmetric tensor; you may be familiar with the $()$ notation for symmetrization, and S_k denotes the set of permutations of k objects. This will be a symmetric tensor, in the sense that, for example you dealt with a symmetric energy-momentum tensor on the previous sheets.

Similarly, there is another corresponding tensor

$$T_{\mu_1, \dots, \mu_k}^{Asym} := T_{[\mu_1, \dots, \mu_k]} = \frac{1}{k!} \sum_{\sigma \in S_K} sgn(\sigma) T_{\mu_{\sigma_1}, \dots, \mu_{\sigma_k}} \quad (4)$$

where again $[\]$ denotes antisymmetrization, implemented by the series on the right. $sgn(\sigma)$ is the sign, i.e. even/odd, of a specific permutation σ . This will be an antisymmetric tensor in the sense that the electromagnetic field strength tensor $F_{\mu\nu}$ was.

- a) Let us specialize to the case $k = 2$, show that $T = T^{sym} + T^{Asym}$. (1 point)

This decomposition is often helpful, to see why, show that for two sets of tensors $X, Y \in T^{0,2}, T^{2,0}$. (1 point)

- b) $X^{sym} Y^{Asym} = 0$.

c) $XY = X^{sym}Y^{sym} + X^{Asym}Y^{Asym}$

Let us now move on to $k = 3$.

d) can you decompose a tensor as $T^{Asym} + T^{sym}$? *(2 points)*