Exercises General Relativity and Cosmology

Priv.-Doz. Stefan Förste, Cesar Fierro

Hand in: 8.5.2017

-Homework-

1 Matter content of the universe as a perfect fluid (10 points)

For a system of N discrete point particles the energy-momentum tensor takes the form

$$T_{\mu\nu} = \sum_{a=1}^{N} \frac{p_{\mu}^{(a)} p_{\nu}^{(a)}}{p^{0(a)}} \delta^{(3)}(\vec{x} - \vec{x^{(a)}}(\sigma_a)), \tag{1}$$

where the index a labels an a-th particle.

a) Show that, for a dense collection of particles with isotropically distributed velocities, we can smooth over the individual particle worldlines to obtain (2 points)

$$T^{\mu\nu} = \begin{pmatrix} \rho & 0 & 0 & 0\\ 0 & p & 0 & 0\\ 0 & 0 & p & 0\\ 0 & 0 & 0 & p \end{pmatrix}$$
(2)

Hint: What would be the definition of p and ρ in this case?.

Note that $T^{\mu\nu}$ and be understood as the stress-energy tensor of a perfect fluid with mass density ρ and pressure p.

b) Show that $T^{\mu\nu}$ can also be written as (2 points)

$$T^{\mu\nu} = (\rho + p)U^{\mu}U^{\nu} + p\eta^{\mu\nu}, \tag{3}$$

where U^{μ} can be tought as the components of the four-velocity of the fluid.

c) From the nonrelativistic limit of the conservation of the energy-momentum tensor, $\partial_{\mu}T^{\mu\nu} = 0$, deduce the Euler's equations (3 points)

$$\partial_t \rho + \nabla \cdot (\rho \vec{v}) = 0, \tag{4}$$

$$\rho \Big[\partial_t \vec{v} + (\vec{v} \cdot \nabla) \vec{v} \Big] = -\nabla p.$$
(5)

Hint: The nonrelativistic limit is given by $U^{\mu} = (1, v^{i}), |v^{i}| \ll 1, p \ll \rho$. Project the equation into pieces along and orthogonal to the four-velocity by contraction with U_{ν} and $P^{\rho}_{\nu} = \delta^{\rho}_{\nu} + U^{\rho}U_{\nu}$.

- d) For the nonrelativistic limit argue that the 'fluid' follows the ideal gas law (2 points).
- e) Show that in the highly relativistic limit, the trace T^{μ}_{μ} gives the equation of state (1 point).

$$p = \frac{1}{3}\rho \tag{6}$$

2 Angular momentum and Noether currents (8 points)

Consider a classical field theory of a scalar field $\phi(x)$ defined over the Minkowski space $\mathbb{R}^{1,3}$ with an action given by

$$S = \int_{\mathbb{R}^{1,3}} \mathrm{d}^4 x \mathcal{L}(\phi, \partial_\mu \phi), \tag{7}$$

which is invariant under Lorentz transformations $x^{\mu} \mapsto \Lambda^{\mu}_{\nu} x^{\nu}$.

a) Show that the parameter of an infinitesimal transformation (1 point)

$$x^{\mu} \mapsto (\delta^{\mu}_{\nu} + \omega^{\mu}_{\nu}) x^{\nu}, \tag{8}$$

has to be antisymmetric in its upper indices, i.e. $\omega^{\mu\nu} + \omega^{\nu\mu} = 0$.

b) Show, that the conserved currents of such an infinitesimal transformation are given by (5 points)

$$\mathcal{M}^{\mu\nu\rho} = x^{\nu}T^{\mu\rho} - x^{\rho}T^{\mu\nu},\tag{9}$$

where $T^{\mu\nu}$ is the canonical energy momentum tensor.

c) What are the conserved charges? Show that (2 points)

$$\frac{d}{dt} \int_{\mathbb{R}^3} \mathrm{d}^3 x x^i T^{00} = \text{const.}$$
(10)

3 Tensor gymnastics (12 points)

Let A be a real valued $n \times n$ dimensional matrix. For the later content of the lecture, the following identity will be useful

$$\det A = \exp\left(\operatorname{tr}\log A\right). \tag{11}$$

(6 points)

a) Prove (11).

In special relativity, a pseudotensor is an object that transform under Lorentz transformations, $x^{\mu'} \mapsto \Lambda^{\mu'}_{\mu} x^{\mu}$, just like a tensor up to a determinant factor.

$$\widetilde{T}^{\mu_1'\cdots\mu_k'}_{\nu_1'\cdots\nu_l'} = (\det\Lambda)\Lambda^{\mu_1'}_{\mu_1}\cdots\Lambda^{\mu_k'}_{\mu_k}\Lambda^{\nu_1}_{\nu_1'}\cdots\Lambda^{\nu_l}_{\nu_l'}\widetilde{T}^{\mu_1\cdots\mu_k}_{\nu_1\cdots\nu_l}.$$
(12)

In particular, the Levi-Civita symbol $\tilde{\epsilon}_{\mu_1\cdots\mu_n}$ is an example of pseudotensors, where

$$\widetilde{\epsilon}_{\mu_1\cdots\mu_n} = \begin{cases}
+1 \text{ if } \mu_1\cdots\mu_n \text{ is an even permutation} \\
-1 \text{ if } \mu_1\cdots\mu_n \text{ is an odd permutation} \\
0 \text{ otherwise}
\end{cases}$$
(13)

b) Show that under a coordinate transformation in special relativity, $\tilde{\epsilon}_{\mu'_1\cdots\mu'_n} = \tilde{\epsilon}_{\mu_1\cdots\mu_n}$. (2 points)

Consider adding to the Lagrangian for electromagnetism of Homework 1 an additional term of the form

$$\widetilde{\mathcal{L}}_{EM} = \widetilde{\epsilon}_{\mu\nu\rho\sigma} F^{\mu\nu} F^{\rho\sigma}.$$
(14)

- c) Express $\widetilde{\mathcal{L}}_{EM}$ in terms of the 3-vector fields \vec{E} and \vec{B} . (2 points)
- d) Show that including $\widetilde{\mathcal{L}}_{EM}$ does not affect Maxwell's equations. Can you think of a deep reason for this? (2 points)