Exercises on Advanced Topics in String Theory

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EXERCISES SHEET 3

3.1 Supersymmetry on the worldsheet

 $(10 \ points)$

Consider the action

$$S = -\frac{1}{2\pi} \int d^2\sigma \left(\partial_\alpha X_\mu \partial^\alpha X^\mu + \overline{\psi}^\mu \rho^\alpha \partial_\alpha \psi_\mu \right) \tag{1}$$

for the superstring. Here ρ^0 and ρ^1 form a representation for the two dimensional Clifford algebra and $\overline{\psi} = \psi^{\dagger}\beta = \psi^{T}\beta$ with $\beta = \rho^0$ is the Dirac conjugate of ψ . $\psi^{\mu} = \begin{pmatrix} \psi^{\mu}_{+} \\ \psi^{\mu}_{-} \end{pmatrix}$ is a two dimensional spinor (it transforms under the spinorial representation of the two dimensional Lorentz group). We choose the components of the spinor to be real i.e. $\psi^{\mu*}_{+} = \psi^{\mu}_{+}$ and $\psi^{\mu*}_{-} = \psi^{\mu}_{-}$ since this is possible in two dimensions.

- (a) Check that $\rho^0 = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$ and $\rho^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$ satisfies the Clifford algebra and evaluate the signature of the worldsheet. (2 points)
- (b) Show that the superstring action can be rewritten as

$$S = \frac{1}{\pi} \int d^2 \sigma \left(2\partial_+ X \partial_- X + i\psi_- \partial_+ \psi_- + i\psi_+ \partial_- \psi_+ \right).$$
(2)

- (2 points)
- (c) Check that (1) is invariant under the global supersymmetry variations given by

$$\delta X^{\mu} = \overline{\epsilon} \psi^{\mu}, \quad \delta \psi^{\mu} = \rho^{\alpha} \partial_{\alpha} X^{\mu} \epsilon, \tag{3}$$

where ϵ is a Majorana spinor.

(d) Show that these induce a conserved supercurrent

$$j_{+} = \psi^{\mu}_{+} \partial_{+} X_{\mu}, \quad j_{-} = \psi^{\mu}_{-} \partial_{-} X_{\mu}.$$
 (4)

(2 points)

(2 points)

(e) Show that the non-zero elements of the energy momentum tensor are given by

$$T_{++} = \partial_{+} X_{\mu} \partial_{+} X^{\mu} + \frac{i}{2} \psi^{\mu}_{+} \partial_{+} \psi_{+\mu}, \quad T_{--} = \partial_{-} X_{\mu} \partial_{-} X^{\mu} + \frac{i}{2} \psi^{\mu}_{-} \partial_{-} \psi_{-\mu}.$$
(5)

 $(2 \ points)$

(10 points)

3.2 Spinors in various dimensions

Let us define the Clifford algebra in d dimensions by

$$\{\gamma^{\mu}, \gamma^{\nu}\} = 2\eta^{\mu\nu},\tag{6}$$

where $\mu, \nu = 0, 1, ..., d - 1$. It holds that $\gamma^{\mu \dagger} = \gamma^{\mu}$ for Euclidean signature, whereas $\gamma^{\mu \dagger} = \gamma^0 \gamma^{\mu} \gamma^0$ for Minkowski space. The 2^{d+1} matrices $\pm 1, \pm \gamma^{\mu}, \pm \gamma^{\mu\nu}, ...$ generate a finite group. Schur's Lemma states that an operator which commutes with all elements of a representation must be a multiple of the unitary element. In addition we introduce the anti-symmetrized products

$$\gamma^{\mu_1\dots\mu_p} = \frac{1}{p!} \left(\gamma^{\mu_1} \gamma^{\mu_2} \dots \gamma^{\mu_p} \pm \text{permutations} \right)$$
(7)

and for even dimensions d = 2n the chirality operator

$$\gamma_{d+1} = \alpha \gamma^0 \gamma^1 \dots \gamma^{d-1}.$$
(8)

The charge conjugation matrix is defined via

$$(\gamma^{\mu})^{\mathrm{T}} = \mp C_{\pm} \gamma^{\mu} C_{\pm}^{-1}. \tag{9}$$

One can show that both C_{\pm}^{T} exist in even dimensions and at least one of them exists in odd dimensions. In particular one has for d = 2n

$$C_{\pm}^{\rm T} = (-1)^{\frac{1}{2}n(n\pm1)} C_{\pm} \tag{10}$$

and for d = 2n + 1

$$C = \begin{cases} C_+ & \text{for } n \text{ odd} \\ C_- & \text{for } n \text{ even} \end{cases}.$$
 (11)

- (a) Determine α s.t $\gamma_d^2 = 1$ holds and show that $\gamma_{d+1}^{\dagger} = \gamma_{d+1}$. (1 point)
- (b) Using Schur's Lemma show that $C^{\mathrm{T}} = \pm C^{\mathrm{T}}$. (1 point)
- (c) Show that the matrices $T^{\mu\nu} = -\frac{i}{2}\gamma^{\mu\nu}$ satisfy

$$[T^{\mu\nu}, T^{\rho\sigma}] = i \left(\eta^{\mu\rho} T^{\nu\sigma} + \eta^{\nu\sigma} T^{\mu\rho} + \eta^{\mu\sigma} T^{\nu\rho} - \eta^{\nu\rho} T^{\mu\sigma}\right).$$
(12)

(1 point)

A representation for which there is a matrix R s.t.

$$-(T^{\mu\nu})^* = RT^{\mu\nu}R^{-1} \tag{13}$$

is called a (pseudo-) real and complex otherwise. In particular one can show that $R^{\rm T} = \pm R$. A representation with a positive sign is called real, whereas the representation with the minus sign is called pseudoreal.

(d) Show that in the even-dimensional Euclidean case d = 2n

$$(T_{\pm}^{\mu\nu})^* = \begin{cases} -(C_{\pm})T_{\mp}^{\mu\nu}(C_{\pm})^{-1} & \text{for } n \text{ odd} \\ -(C_{\pm})T_{\pm}^{\mu\nu}(C_{\pm})^{-1} & \text{for } n \text{ even} \end{cases},$$
(14)

where $T_{\pm}^{\mu\nu} = T^{\mu\nu}\frac{1}{2}(1 \pm \gamma_{d+1})$ are the generators associated to the respective chiral subspaces. (1 point)

- (e) Evaluate for which dimensions the representations are real, pseudoreal and complex. You should find that the result only depends on the dimension mod 4. (1 point)
- (f) We define in Euclidean signature

$$b_i^{\pm} = \frac{1}{2} (\gamma^{2i} \pm i\gamma^{2i+1}). \tag{15}$$

Show that

$$b_i^{\pm} = b_i^{\pm\dagger}, \quad \{b_i^{\pm}, b_j^{\mp}\} = \delta_{ij}, \quad \{b_i^{+}, b_j^{+}\} = \{b_i^{-}, b_j^{-}\} = 0.$$
(16)

(2 points)

- (g) Let the highest weight state $|\Omega\rangle$ be defined by $b^i |\Omega\rangle = 0$ such that $|\Omega\rangle = |\frac{1}{2}, ..., \frac{1}{2}\rangle$. All other states are given by $|\pm \frac{1}{2}, ..., \pm \frac{1}{2}\rangle$. Show that this representation is reducible and decomposes into irreducible representations given by positive and negative chirality spinors respectively. (2 points)
- (h) Show that d = 8 is special in the sense that the spinorial representations have the same dimension as the vector representation. There is a symmetry relation these representations, called triality symmetry. Can you guess this symmetry by inspecting the Dynkin diagram of SO(8)? (1 point)