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## Exercises on General Relativity and Cosmology

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<http://www.th.physik.uni-bonn.de/people/forste/exercises/ss2013/gr>

### –HOME EXERCISES–

#### H 3.1 Some concrete tensor algebra

(4 points)

Given the components of a  $(2,0)$ -tensor  $X$  as well as the components of a vector  $V$ ,

$$(X^{\mu\nu}) = \begin{pmatrix} 2 & 0 & 1 & -1 \\ -1 & 0 & 3 & 2 \\ -1 & 1 & 0 & 0 \\ -2 & 1 & 1 & -2 \end{pmatrix}, \quad (V^\mu) = \begin{pmatrix} -1 \\ 2 \\ 0 \\ -2 \end{pmatrix},$$

compute

- (a)  $X^\mu{}_\nu$ ,
- (b)  $X_\mu{}^\nu$ ,
- (c)  $X^{(\mu\nu)}$ ,
- (d)  $X^{[\mu\nu]}$ ,
- (e)  $X^\mu{}_\mu$ ,
- (f)  $V^\mu V_\mu$ ,
- (g)  $V_\mu X^{\mu\nu}$ .

#### H 3.2 Energy-momentum tensor

(16 points)

Given a collection of charged particles with positions  $\vec{x}_n(t)$  and charges  $e_n$  we define electric charge density  $\rho$  and current density  $\vec{j}$  as

$$\rho(\vec{x}, t) = \sum_n e_n \delta^3(\vec{x} - \vec{x}_n(t)), \quad \vec{j}(\vec{x}, t) = \sum_n e_n \dot{\vec{x}}_n(t) \delta^3(\vec{x} - \vec{x}_n(t)).$$

Similarly, we define the charge density for the four-momentum  $p^\mu$ , the *energy-momentum tensor*, as

$$T^{\mu\nu} = \sum_n p_n^\mu(t) \dot{x}_n^\nu(t) \delta^3(\vec{x} - \vec{x}_n(t)).$$

- (a) Check that the  $T^{\mu\nu}$  transform as components of a  $(2,0)$ -tensor. (1 point)

- (b) Show that the energy-momentum tensor is only conserved up to a force density  $G^\mu$  which vanishes for free particles

$$\partial_\nu T^{\mu\nu} = G^\mu .$$

(2 points)

- (c) Check that for the electromagnetic force from exercise H2.2(c),

$$f^\mu \equiv \frac{dp^\mu}{d\tau} = e F^\mu{}_\nu \frac{dx^\nu}{d\tau} ,$$

$G$  is given by

$$G^\mu = F^\mu{}_\nu J^\nu .$$

(1 point)

- (d) To obtain a conserved energy-momentum tensor, we have to include the contribution of the electromagnetic field itself

$$T_{\text{em}}^{\mu\nu} = F^\mu{}_\rho F^{\nu\rho} - \frac{1}{4} \eta^{\mu\nu} F^{\rho\sigma} F_{\rho\sigma} .$$

Show that  $\partial_\nu T_{\text{em}}^{\mu\nu}$  cancels  $G^\mu$  introduced in point (b). Thus,  $T_{\text{tot.}}^{\mu\nu} = T^{\mu\nu} + T_{\text{em}}^{\mu\nu}$  is conserved. (2 points)

- (e) Show that the total momentum

$$P^\mu = \int d^3x T_{\text{tot.}}^{\mu 0}(\vec{x}, t)$$

is a conserved quantity.

(2 points)

Now we want to consider the energy-momentum tensor of a *perfect fluid*. A comoving observer will, by definition, see his surroundings as isotropic. In this frame the energy-momentum tensor is given by

$$\left( \tilde{T}^{\mu\nu} \right) = \begin{pmatrix} \rho & 0 & 0 & 0 \\ 0 & p & 0 & 0 \\ 0 & 0 & p & 0 \\ 0 & 0 & 0 & p \end{pmatrix} ,$$

where  $\rho$  is the density and  $p$  the pressure of the fluid.

- (f) Calculate the components of energy-momentum tensor  $T^{\mu\nu}$  for an observer at rest. Assume the comoving observer's velocity to be  $\vec{v}$ . (3 points)

- (g) Show that  $T^{\mu\nu}$  can also be written as

$$T^{\mu\nu} = (p + \rho) U^\mu U^\nu + p \eta^{\mu\nu} ,$$

where  $U^\mu$  are the components of the four-velocity of the fluid.

(2 points)

- (h) From the nonrelativistic limit of the conservation of the energy momentum tensor,  $\partial_\mu T^{\mu\nu}$ , deduce *Euler's equations*

$$\begin{aligned}\partial_t \rho + \vec{\nabla} \cdot (\rho \vec{v}) &= 0, \\ \rho \left[ \partial_t \vec{v} + (\vec{v} \cdot \vec{\nabla}) \vec{v} \right] &= -\vec{\nabla} p.\end{aligned}$$

*Hint: The nonrelativistic limit is given by  $(U^\mu) = (1, v^i)$ ,  $|v^i| \ll 1$ ,  $p \ll \rho$ . Project the equation into pieces along and orthogonal to the four-velocity by contraction with  $U_\nu$  and  $P^\sigma{}_\nu = \delta^\sigma_\nu + U^\sigma U_\nu$  respectively.* (3 points)