

## Superstring Theory

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<http://www.th.physik.uni-bonn.de/people/forste/exercises/strings19>

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–HOMEWORKS–

### 3.1 Oscillator expansions for the classical string

In exercise 2.2 we have derived the equations of motion for the classical string in conformal gauge, i.e.  $(\partial_\sigma^2 - \partial_\tau^2) X^\mu$ . The vanishing of the surface term followed after imposing any of the three following boundary conditions

- i) Closed String:  $X^\mu(\sigma, \tau) = X^\mu(\sigma + l, \tau)$ ,
- ii) Open String (Neumann):  $X'_\mu(\sigma, \tau) = 0$  for  $\sigma = 0, l$ ,
- iii) Open String (Dirichlet):  $X^\mu|_{\sigma=0} = X_0^\mu = \text{const.}$  and  $X^\mu|_{\sigma=l} = X_l^\mu = \text{const.}$ .

We make a product ansatz for a solution of the equations of motion to any of the  $d$  spacetime coordinates  $X^\mu$  with  $\mu = 0, \dots, d-1$  given by

$$X^\mu(\sigma, \tau) = f(\sigma)g(\tau).$$

We start our discussion with the *closed string*.

- a) Show that the functions  $f(\sigma)$  and  $g(\tau)$  have to satisfy

$$\frac{\partial^2 f(\sigma)}{\partial \sigma^2} = k f(\sigma), \quad \frac{\partial^2 g(\tau)}{\partial \tau^2} = k g(\tau), \quad k = -\frac{4m^2\pi^2}{l^2}, \quad m \in \mathbb{Z}.$$

Solve these differential equations for  $m = 0$  as well as for  $m \in \mathbb{Z} \setminus \{0\}$ . Out of this write the full solution  $X^\mu(\sigma, \tau)$ . (2 Points)

- b) Argue that the general solution for a *closed string* splits into left-movers and right-movers

$$X^\mu(\sigma, \tau) = X_L^\mu(\sigma^+) + X_R^\mu(\sigma^-).$$

With a convenient normalization (inspired from the analysis in previous part) the left-movers and the right-movers are respectively given by

$$X_L^\mu(\sigma^+) = \frac{1}{2}(x^\mu - c^\mu) + \frac{\pi\alpha'}{l}p^\mu\sigma^+ + i\sqrt{\frac{\alpha'}{2}} \sum_{n \neq 0} \frac{1}{n} \tilde{\alpha}_n^\mu e^{-\frac{2\pi}{l}in\sigma^+},$$

$$X_R^\mu(\sigma^-) = \frac{1}{2}(x^\mu - c^\mu) + \frac{\pi\alpha'}{l}p^\mu\sigma^- + i\sqrt{\frac{\alpha'}{2}} \sum_{n \neq 0} \frac{1}{n} \alpha_n^\mu e^{-\frac{2\pi}{l}in\sigma^-}.$$

In these equations  $n \in \mathbb{Z} \setminus \{0\}$ ,  $\alpha_n^\mu$  and  $\tilde{\alpha}_n^\mu$  are arbitrary Fourier modes with  $\alpha_n^\mu, \tilde{\alpha}_n^\mu$  positive-frequency modes for  $n < 0$  and  $\alpha_n^\mu, \tilde{\alpha}_n^\mu$  negative-frequency modes for  $n > 0$  and  $c^\mu$  is an arbitrary parameter which can be set to zero if the zero-mode part of the expansions is left-right symmetric. (1 Point)

- c) Take the solution from part a) and give conditions which have to be imposed on  $x^\mu, p^\mu, \alpha_n^\mu$  and  $\tilde{\alpha}_n^\mu$  in order to make  $X^\mu$  real. (1 Point)
- d) We define  $\alpha_0^\mu = \tilde{\alpha}_0^\mu = \sqrt{\frac{\alpha'}{2}} p^\mu$ . Compute the following quantities in terms of the oscillator modes appearing in the left-movers  $X_L^\mu(\sigma^+)$  and the right-movers  $X_R^\mu(\sigma^-)$
- i) total spacetime momentum:  $P^\mu = \frac{1}{2\pi\alpha'} \int_0^l d\sigma \dot{X}^\mu$ ,
  - ii) center of mass position:  $q^\mu = \frac{1}{l} \int_0^l d\sigma X^\mu$ ,
  - iii) total angular momentum:  $J^{\mu\nu} = \frac{1}{2\pi\alpha'} \int_0^l d\sigma \left( X^\mu \dot{X}^\nu - X^\nu \dot{X}^\mu \right)$ .

Furthermore, explain the physical interpretation of  $p^\mu$  and  $x^\mu$  for the string. (3 Points)

Now we want to obtain the oscillator expansions for *open strings*.

- e) Argue that open string solutions have only one set of oscillator modes. (1 Point)
- f) Show that the mode expansion for *open strings* with *Neumann* boundary conditions at both ends of the string, i.e. for  $\sigma = 0$  and  $\sigma = l$ , is given by

$$X^\mu(\sigma, \tau) = x^\mu + \frac{2\pi\alpha'}{l} p^\mu \tau + i\sqrt{2\alpha'} \sum_{n \neq 0} \frac{1}{n} \alpha_n^\mu e^{-i\frac{\pi}{l} n \tau} \cos\left(\frac{n\pi\sigma}{l}\right). \quad (1 \text{ Point})$$

- g) Show that the mode expansion for *open strings* with *Dirichlet* boundary conditions at both ends of the string, i.e. for  $\sigma = 0$  and  $\sigma = l$ , is given by

$$X^\mu(\sigma, \tau) = x_0^\mu + \frac{1}{l} (x_l^\mu - x_0^\mu) \sigma + \sqrt{2\alpha'} \sum_{n \neq 0} \frac{1}{n} \alpha_n^\mu e^{-i\frac{\pi}{l} n \tau} \sin\left(\frac{n\pi\sigma}{l}\right). \quad (1 \text{ Point})$$

### 3.2 Poisson brackets for the classical closed string

In this exercise we want to find the Poisson brackets for the oscillator modes  $\alpha_n^\mu, \tilde{\alpha}_n^\mu$ . Recall the equal time Poisson brackets

$$\begin{aligned} \{X^\mu(\sigma, \tau), X^\nu(\sigma', \tau)\}_{\text{P.B.}} &= \{\dot{X}^\mu(\sigma, \tau), \dot{X}^\nu(\sigma', \tau)\}_{\text{P.B.}} = 0, \\ \{X^\mu(\sigma, \tau), \dot{X}^\nu(\sigma', \tau)\}_{\text{P.B.}} &= \frac{1}{T} \eta^{\mu\nu} \delta(\sigma - \sigma'). \end{aligned}$$

Use the oscillator expansion for the closed string from the previous exercise 3.1 to show that the Poisson brackets for the modes are given by

$$\begin{aligned} \{\alpha_n^\mu, \alpha_m^\nu\}_{\text{P.B.}} &= \{\tilde{\alpha}_n^\mu, \tilde{\alpha}_m^\nu\}_{\text{P.B.}} = -in\eta^{\mu\nu} \delta_{m+n,0}, & \{\alpha_n^\mu, \tilde{\alpha}_m^\nu\}_{\text{P.B.}} &= 0 \\ \{x^\mu, x^\nu\}_{\text{P.B.}} &= \{p^\mu, p^\nu\}_{\text{P.B.}} = 0, & \{x^\mu, p^\nu\}_{\text{P.B.}} &= \eta^{\mu\nu}. \end{aligned}$$

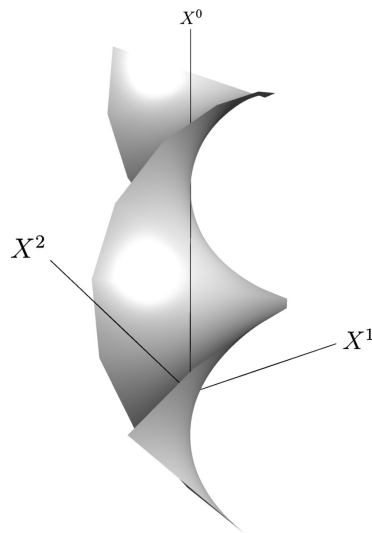
Hint: Express  $\alpha_n^\mu$  and  $\tilde{\alpha}_n^\mu$  as linear combinations of  $X^\mu(\sigma, \tau)$  and  $\dot{X}^\mu(\sigma, \tau)$  for fixed  $\tau$ . Then you can use the equal time Poisson brackets. (4 Points)

### 3.3 A spaghetti stick as solution to the string equations of motion

We consider a with constant angular velocity rotating spaghetti stick in the  $X^1$ - $X^2$ -plane parametrized by

$$X^0 = A\tau, \quad X^1 = A \cos \tau \cos \sigma, \quad X^2 = A \sin \tau \cos \sigma, \quad X^i = 0 \quad \text{for } i = 3, \dots, d-1.$$

The worldsheet of this configuration looks like



Throughout this exercise we want to show that this configuration is an open string solution and compute its characteristics.

- Verify that the rotating spaghetti stick is indeed an open string solution. (1 Point)
- Determine the boundary conditions of this configuration and the speed of the string's endpoints. (1 Point)
- Compute the energy  $M := P^0$  of the spaghetti stick. (1 Point)
- Compute the angular momentum  $J := |J^{12}|$ . (1 Point)
- Compute the quantity  $\frac{J}{M^2}$ . (1 Point)