

Übungen zu Theoretische Physik IV

Priv.-Doz. Dr. Stefan Förste

<http://www.th.physik.uni-bonn.de/people/forste/exercises/ws1213/tp4>

–IN-CLASS EXERCISES–

A 2.1 Saddle-point method

In the context of statistical physics it is often necessary to solve integrals of the form

$$I = \lim_{N \rightarrow \infty} \int_a^b e^{Nf(x)} dx.$$

If $f(x)$ is an analytic function on the interval $[a, b]$ and has a global minimum at $x_0 \in (a, b)$ then

$$I = \lim_{N \rightarrow \infty} e^{Nf(x_0)} \sqrt{\frac{2\pi}{N|f''(x_0)|}}, \quad \text{where } f''(x_0) \equiv \left. \frac{\partial^2 f(x)}{\partial x^2} \right|_{x=x_0}.$$

Show *Stirling's formula*

$$N! \rightarrow \sqrt{2\pi N} N^N e^{-N} \quad \text{as } N \rightarrow \infty$$

using the saddle-point method.

Hint: Use the integral representation of the gamma function $N! = \Gamma(N+1) = \int_0^\infty x^N e^{-x} dx$.

A 2.2 Ensemble of quantum mechanical harmonic oscillators

Consider a system of N distinguishable, non-interacting, quantum mechanical, harmonic oscillators with equal angular velocity ω . States of the complete system are then given by the individual oscillator states

$$|n_1, n_2, \dots, n_N\rangle = |n_1\rangle \otimes |n_2\rangle \otimes \dots \otimes |n_N\rangle.$$

We write shorthand

$$a_i \equiv \mathbb{1}^{\otimes(i-1)} \otimes a \otimes \mathbb{1}^{\otimes(N-i)} = \mathbb{1} \otimes \dots \otimes \underset{\substack{\uparrow \\ \text{ite Stelle}}}{a} \otimes \mathbb{1} \otimes \dots \otimes \mathbb{1}$$

for the *lowering operator* of the i th oscillator (with similar expressions for a_j^\dagger , N_j , H_j). The Hamilton operator of the system is given by

$$H = \sum_{j=1}^N \hbar\omega \left(a_j^\dagger a_j + \frac{1}{2} \right).$$

First, consider the case $N = 3$ with total energy $E = \frac{9}{2}\hbar\omega$ of the system.

- (a) How many different states realize this value of the energy?
- (b) What is the probability $p(\epsilon)$ for a given oscillator to have the energy ϵ ?

Now we want to determine the number of states for a given energy E in the limit of a very big number of oscillators $N \gg 1$. In general, it is given by

$$\Omega(E) \equiv \text{Sp } \delta(E - H).$$

- (c) What is $\Omega(E)$ for the given system?
- (d) Show, that

$$\Omega(E) = \int \frac{dk}{2\pi} e^{ikE} \left(\frac{e^{-ik\hbar\omega/2}}{1 - e^{-ik\hbar\omega}} \right)^N$$

and further, that

$$\Omega(E) = \int \frac{dk}{2\pi} e^{N(ik(E/N) - \log(2i \sin(k\hbar\omega/2)))}.$$

- (e) This integral can be computed using the saddle-point method. Show, that $\Omega(E)$ is given by

$$\Omega(E) = \exp \left\{ N \left[\frac{\frac{E}{N} + \frac{1}{2}\hbar\omega}{\hbar\omega} \log \frac{\frac{E}{N} + \frac{1}{2}\hbar\omega}{\hbar\omega} - \frac{\frac{E}{N} - \frac{1}{2}\hbar\omega}{\hbar\omega} \log \frac{\frac{E}{N} - \frac{1}{2}\hbar\omega}{\hbar\omega} \right] \right\}.$$

–HOMEWORK–

H 2.1 Spin precession of a spin-1/2 particle

(5+5=10) Points

The Hamilton operator of a spin-1/2 particle in a homogeneous magnetic field B is given by

$$H = -\frac{\gamma}{2}\hbar \sum_{j=1}^3 \sigma_j B_j,$$

where γ is the gyromagnetic constant and σ_i ($i = 1, 2, 3$) are the Pauli matrices. The time variation of the polarisation is given by

$$\frac{\partial P_i}{\partial t} = \frac{\partial \langle \sigma_i \rangle}{\partial t}.$$

Here, σ_i is to be understood as an operator in the spin-1/2 representation of the rotation group.

- (a) Express the time variation of the polarisation in terms of the density matrix $\rho(t)$ and show

$$i\frac{\partial P_i}{\partial t} = -\frac{\gamma}{2} \sum_j B_j \text{Sp}([\sigma_i, \sigma_j]\rho).$$

- (b) Show the *Bloch equation*

$$\frac{\partial}{\partial t} P = \gamma (P \times B).$$

$$\text{Hint: } [\sigma_i, \sigma_j] = 2i \sum_k \sigma_k \epsilon_{ijk}.$$

H 2.2 Spin ensemble

(4+2+2+2=10) Points

Consider a system of N ($N \gg 1$) non-interacting spin-1/2 particles in a constant magnetic field B . Each of the particles has a magnetic moment μ , which can be aligned parallelly or anti-parallelly to the magnetic field. Let n_1 (n_2) be the number of magnetic moments aligned parallelly (anti-parallelly) to the magnetic field. The energy of the system is then given by $E = -(n_1 - n_2)\mu B$.

- (a) Show, that the number of states having an energy between E and $E + \delta E$ is given approximately by

$$\omega(E, \delta E) = \frac{N!}{\left(\frac{N}{2} - \frac{E}{2\mu B}\right)! \left(\frac{N}{2} + \frac{E}{2\mu B}\right)!} \frac{\delta E}{2\mu B},$$

where $E \gg \delta E \gg \mu B$.

Hint: What is the energy difference between two energy levels?

- (b) Use Stirling's formula, as it was derived in exercise A 2.1, to find an approximation to $\ln \omega(E, \delta E)$.
- (c) Interpret the function $\ln f(n_1) \equiv \ln \left(\frac{N!}{n_1!(N-n_1)!} \right)$ as a continuous function of n_1 . Its Taylor expansion around the maximum $n_{1,\max}$ up to second order is then given by

$$\ln f(n_1) = \ln f(n_{1,\max}) + \frac{1}{2}(n_1 - n_{1,\max})^2 \left. \frac{\partial^2 \ln f(n_1)}{\partial n_1^2} \right|_{n_1=n_{1,\max}}.$$

Using the approximation $\ln n! \approx n \ln n - n$, calculate the maximum $n_{1,\max}$ as well as $\left. \frac{\partial^2 \ln f(n_1)}{\partial n_1^2} \right|_{n_1=n_{1,\max}}$.

- (d) Use the exponential of the Taylor expansion of $\ln f(n_1)$ to show that $\omega(E, \delta E)$ is given approximately by the Gaussian distribution

$$\omega(E, \delta E) = 2^N \sqrt{\frac{2}{\pi N}} \frac{\delta E}{2\mu B} \exp \left[-\frac{2}{N} \left(\frac{E}{2\mu B} \right)^2 \right].$$