

Condensed Matter Theory I — WS05/06

Exercise 3

(Please return your solutions before 22.11., 13:00 h)

3.1 Density of states (DOS)

(4 points)

- a) Derive the DOS $N(E)$ for a general dispersion relation ϵ_k in d dimensions, starting with the definition:

$$N(E) := \sum_k \delta(E - \epsilon_k) \quad (1)$$

(*hint*: Use the equation for the group velocity, given in the lecture.)

- b) Calculate $N(E)$ for the 1d chain ($\epsilon_k = -2t\cos(ka)$).

3.2 Green's functions - equation of motion

(8 points)

- a) Derive an equation of motion for an operator A , which does not depend explicitly on time, in the Heisenberg picture.
- b) Determine the time-dependence of $c_{k\sigma}(t)$ and $c_{k\sigma}^\dagger(t)$ for the free electron gas with $H = \sum_{k\sigma} \epsilon_k c_{k\sigma}^\dagger c_{k\sigma}$.
- c) Compute the free retarded Green's function $G_{k\sigma}^{0,R}(t - t')$, using the result of b).
- d) Derive the Fourier transform $G_{k\sigma}^{0,R}(\omega) = \int d\tau G_{k\sigma}^{0,R}(\tau) e^{i(\omega + i\eta)\tau}$ with $|\eta| \rightarrow 0$. What is the sign of η ? Why?

3.3 Equation of motion for the Green's function with local pair interaction

(10 points)

If one considers an electron gas with a local pair interaction between two electrons, the Hamiltonian in position representation in 2nd quantization looks like this:

$$\begin{aligned}
H_0 &= \frac{\hbar^2}{2m} \sum_{\sigma} \int d^3x \nabla \Psi_{\sigma}^{\dagger}(x) \nabla \Psi_{\sigma}(x) \\
V &= \frac{1}{2} \sum_{\sigma\sigma'} \int d^3x d^3x' v(x, x') \Psi_{\sigma}^{\dagger}(x) \Psi_{\sigma'}^{\dagger}(x') \Psi_{\sigma'}(x') \Psi_{\sigma}(x)
\end{aligned} \tag{2}$$

where the pair interaction is given by: $v(x, x') = (2\pi)^3 V_0 \delta(x - x') \cdot \delta_{\sigma' - \sigma}$.

In a real system, the local pair interaction can be explained by a strongly screened Coulomb interaction.

- a) Transform the Hamiltonian (Eqn.2) to momentum space and show, using $\epsilon_k = -\frac{\hbar^2 k^2}{2m}$, that:

$$H = H_0 + V = \sum_{k\sigma} \epsilon_k c_{k\sigma}^{\dagger} c_{k\sigma} + \frac{1}{2} V_0 \sum_{\substack{kpq \\ \sigma\sigma'}} \delta_{\sigma' - \sigma} c_{k+q\sigma}^{\dagger} c_{p-q\sigma'}^{\dagger} c_{p\sigma'} c_{k\sigma} \tag{3}$$

- b) The equation of motion of the retarded Green's function for the Hamiltonian H was shown in the lecture to be:

$$\begin{aligned}
i\hbar \partial_t G_{\kappa\tau}^R(t, t') &= \hbar \delta(t - t') \langle [c_{\kappa\tau}, c_{\kappa\tau}^{\dagger}]_+ \rangle - i\theta(t - t') \langle [[c_{\kappa\tau}, H_0]_-(t), c_{\kappa\tau}^{\dagger}(t')]_+ \rangle \\
&\quad - i\theta(t - t') \langle [[c_{\kappa\tau}, V]_-(t), c_{\kappa\tau}^{\dagger}(t')]_+ \rangle
\end{aligned}$$

Write down the equation of motion for the Hamiltonian of part a) and show that the interaction term depends on the higher Green's function

$$\Gamma_{pq\kappa}^{-\tau\tau} = -i\theta(t - t') \langle [(c_{p+q-\tau}^{\dagger} c_{p-\tau} c_{\kappa+q\tau})(t), c_{\kappa\tau}^{\dagger}(t')]_+ \rangle \tag{4}$$

- c) For the following approximation, fluctuations of the operators A and B around their expectation values are neglected, i.e. one considers $A - \langle A \rangle = B - \langle B \rangle = 0$. Thus show that in general the following relation holds: $AB = \langle A \rangle B + \langle B \rangle A - \langle A \rangle \langle B \rangle$.

- d) Use the approximation of part c) with $A = c_{p+q\sigma}^{\dagger} c_{p\sigma}$, $B = c_{\kappa+q\tau}$ and show that the retarded Green's function for the Hamiltonian of Eqn. (2) in momentum space becomes:

$$G_{\kappa\tau}^R(\omega) = \frac{1}{\omega - \epsilon_{\kappa} - V_0 \sum_p \langle n_{p-\tau} \rangle + i\eta}$$

In this case the approximation of part c) is called *mean field approximation*.

- e) Derive an expression for the magnetization $m := \langle n_{\uparrow} \rangle - \langle n_{\downarrow} \rangle$ and the particle number $n := \langle n_{\uparrow} \rangle + \langle n_{\downarrow} \rangle$ and show, that $m = 0$ is always a solution. One can show, that for $1 \leq V_0 N(\epsilon_F)$ there exist (numerical) solutions (Fig. 1) with finite magnetization.

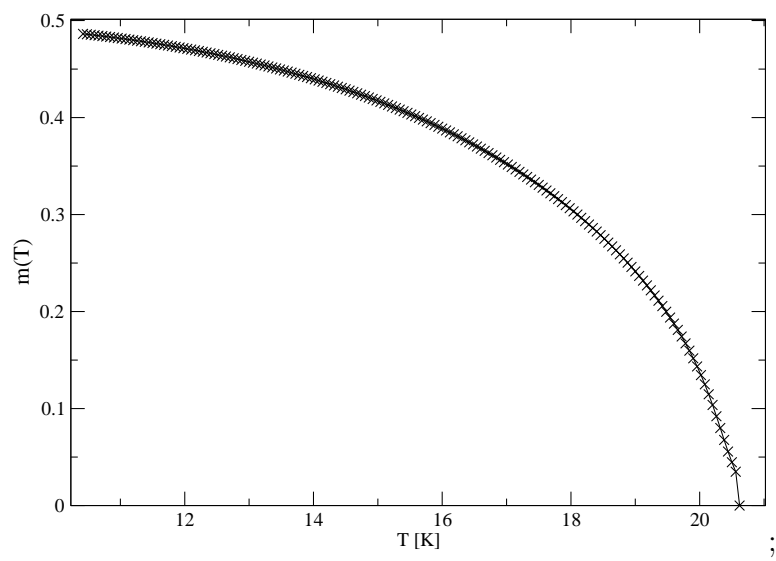


Figure 1: Example for the numerical solution of the magnetization in the mean-field approximation